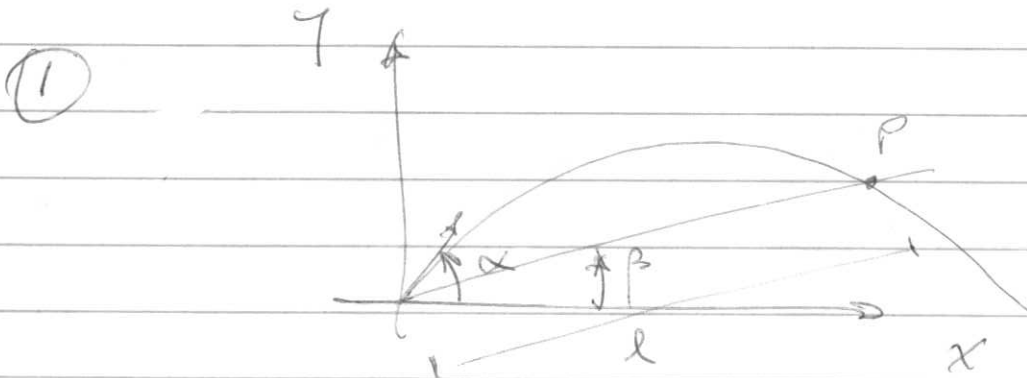


PROBLEMS # 2



The equation of motion is

$$\vec{F} = m\vec{a}$$

The gravitational force is the only applied force; therefore

$$F_x = m\ddot{x} = 0$$

$$F_y = m\ddot{y} = -mg$$

Integrating these equations and using

the initial conditions

$$\dot{x}(t=0) = v_0 \cos \alpha$$

$$\dot{y}(t=0) = v_0 \sin \alpha$$

we find

$$\dot{x}(t) = v_0 \cos \alpha$$

$$\dot{y}(t) = v_0 \sin \alpha - gt$$

So the equations for x and y are

$$x(t) = v_0 t \cos \alpha$$

$$y(t) = v_0 t \sin \alpha - \frac{1}{2} g t^2$$

Suppose it takes a time t_0 to reach

the point P . Then

$$l \cos \beta = v_0 t_0 \cos \alpha$$

$$l \sin \beta = v_0 t_0 \sin \alpha - \frac{1}{2} g t_0^2$$

Eliminating l between these equations

$$\frac{1}{2} g t_0^2 \left[t_0 - \frac{2 v_0 \sin \alpha}{g} + \frac{2 v_0 \cos \alpha \tan \beta}{g} \right] = 0$$

from which

$$t_0 = \frac{2 v_0}{g} (\sin \alpha - \cos \alpha \tan \beta)$$

(2) The static frictional force has the approximate value $f_{\max} = \mu_s N$

$$\vec{F} = \vec{f} + \vec{N}$$

y-direction

$$-f \cos \theta + N = 0$$

x-direction

$$-f_s + f_g \sin \theta = m \ddot{x}$$

The static frictional force f_s will be some value $f_s \leq f_{\max}$ required to keep $\ddot{x} = 0$. However, as the angle

θ of the plane increases, eventually the static frictional force will be unable to keep the block at rest.

At that angle θ' , f becomes

$$f_s(\theta = \theta') = f_{\max} = \mu_s N = \mu_s F_g \cos \theta$$

and

$$m \ddot{x} = F_g \sin \theta - f_{\max}$$

$$m \ddot{x} = F_g \sin \theta - \mu_s F_g \cos \theta$$

$$\ddot{x} = g (\sin \theta - \mu_s \cos \theta)$$

Just before the block starts to slide

the acceleration $\ddot{x} = 0 \Rightarrow$

$$\sin \theta - \mu_s \cos \theta = 0$$

$$\tan \theta = \mu_s = 0.4$$

$$\theta = 22^\circ$$

(ii) the kinetic friction force approximately

$$f_k = \mu_k N = \mu_k F_g \cos \theta$$

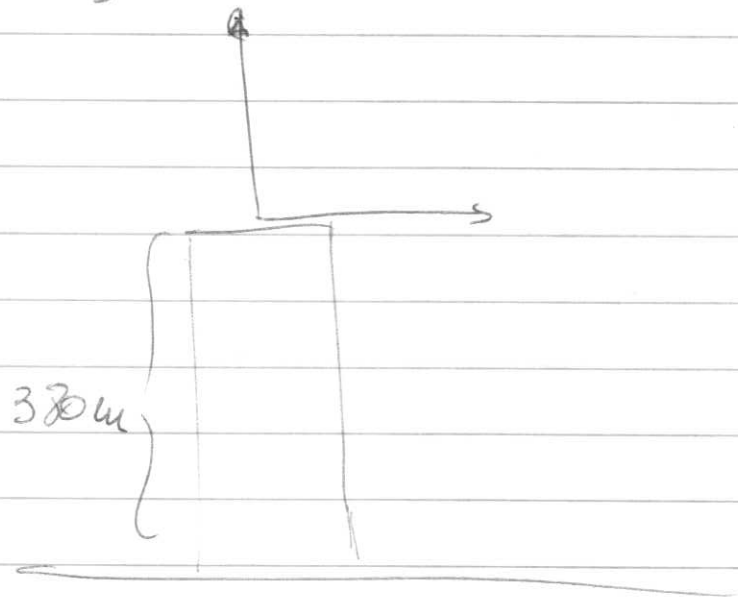
and

$$m \ddot{x} = F_g \sin \theta - f_k = \mu_g (\sin \theta - \mu_k \cos \theta)$$

$$\ddot{x} = g (\sin \theta - \mu_k \cos \theta) = 0.24g$$

(3) Take origin at the top of Empire

State Building



$$m \ddot{y} = -mg$$

$$\frac{dv}{dt} = -g \Rightarrow v = v_0 - gt$$

$v_0 = 0$ from initial condition

$$\frac{dy}{dt} = -gt \Rightarrow y = y_0 - \frac{1}{2}gt^2$$

with $y_0 = 0 \Rightarrow$ from initial condition

$$y_f = 380 \text{ m} \quad t = \sqrt{\frac{2y_0}{g}} = 3.85 \text{ s}, \quad (\ddot{y}) v = 86 \text{ m/s}$$

$$(iii) \frac{dV}{dt} = -g - kV$$

$$\int_0^t dt = \int_0^V \frac{-1}{g + kV} dV$$

$$u = g + kV \quad du = k dV$$

$$\Rightarrow t = \left(\frac{-1}{k} \right) \int_{u_1}^{u_2} du$$

$$-kt = \ln(g + kV) - \ln g$$

$$-kt = \ln \left(\frac{g + kV}{g} \right)$$

$$e^{-kt} = \frac{g + kV}{g}$$

$$\Rightarrow \text{for } t=0 \Rightarrow V=0 \Rightarrow 1=1.$$

$$V = \frac{g}{k} e^{-kt} - \frac{g}{k}$$

$$\frac{dy}{dt} = v = \frac{g e^{-kt}}{k} - \frac{g}{k}$$

$$\Rightarrow \int_0^y dy = \int_0^t \left[\frac{g e^{-kt}}{k} - \frac{g}{k} \right] dt$$

$$\Rightarrow y = \left. -\frac{g e^{-kt}}{k^2} \right|_0^t - \left. \frac{g t}{k} \right|_0^t$$

$$y = -\frac{g e^{-kt}}{k^2} + \frac{g}{k^2} - \frac{g t}{k}$$

$$y=0 \text{ at } t=0$$

(iv)

If we neglect air resistance

$$v = 19.6 \text{ m/s}$$

with the correction of retarding force

$$v = 18.94 \text{ m/s}$$

$$\frac{19.6 - 18.94}{19.60} \times 100\% \approx 3\%$$

(4)

Angular momentum must be conserved during the process. We use the concept of moment of inertia from introductory physics to relate angular momentum L to angular velocity: $L = I\omega$. The initial angular momentum $L_0 = I\omega_0$ must be equal to the angular momentum L (for + mouse) after the mouse jumps on. The velocity of the outside edge is $v = \omega R$.

$$L = I\omega + m v R = \frac{v}{R} (I + m R^2)$$

$$L = L_0 = I\omega_0$$

$$\frac{v}{R} (I + m R^2) = I \frac{v_0}{R}$$

$$\frac{v}{v_0} = \frac{I}{I + m R^2}$$

and

$$\frac{\omega}{\omega_0} = \frac{I}{I + m R^2}$$

(5)⁽²⁾ If here is no air resistance, then conservation of mechanical energy can be used. Subscript 1 represent the glider when it launched, and subscript 2 represent the glider at landing. The landing position is the zero for the potential energy $PE(x=0)$. we have $y_1 = 500 \text{ m}$ and $y_2 = 0$.

$$v_1 = 500 \text{ km/h} \left(\frac{1 \text{ m/s}}{3.6 \text{ km/h}} \right) = 138.9 \text{ m/s}$$

Solve for v_2

$$E_1 = E_2 \Rightarrow \frac{1}{2} m v_1^2 + m g y_1 = \frac{1}{2} m v_2^2 + m g y_2$$

$$v_2 = \sqrt{v_1^2 + 2 g y_1} = 296 \text{ m/s} \left(\frac{3.6 \text{ km/h}}{1 \text{ m/s}} \right)$$

$$= 1067 \text{ km/h} \approx 1.1 \cdot 10^3 \text{ km/h}$$

(b) Now include the work done by the non-conservative force, consider the diagram of the glider



calculate the work done by friction

$$W_{nc} = F_{fr} d \cos 180^\circ = -F_{fr} d = -F_{fr} \frac{3500 \text{ m}}{\sin 10^\circ}$$

Use the same subscript representations as above with y_1, v_1, y_2 as before and

$$v_2 = 200 \text{ km/h} \left(\frac{1 \text{ m/s}}{3.6 \text{ km/h}} \right) = 55.56 \text{ m/s}$$

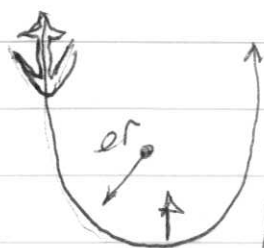
Write the energy conservation equation and solve for the frictional force

$$W_{nc} + E_1 = E_2 \Rightarrow$$

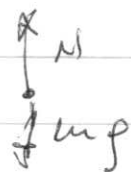
$$F_{fr} d + \frac{1}{2} m v_1^2 + m g y_1 = \frac{1}{2} m v_2^2 + m g y_2$$

$$\Rightarrow F_{fr} = m \frac{(v_1^2 - v_2^2 + 2g y_1)}{2d} \approx 2 \times 10^3 \text{ N}$$

(4)



point of maximum acceleration



(a) From the force diagram we

$$\text{have } (N - mg) \hat{e}_r = \left(\frac{mv^2}{R} \right) \hat{e}_r$$

The acceleration that the pilot feels is

$$\bar{a} = \left(\frac{d}{dt} \right) \hat{e}_r = g + (v^2/R) \hat{e}_r, \text{ which has}$$

maximum magnitude at the bottom of
the manoeuvre

(4) If the acceleration felt by the pilot
must be less than $9g$, then we have

$$R \geq \frac{v^2}{8g} = \frac{(3 \times 340 \text{ m s}^{-1})^2}{8 \times 9.8 \text{ m/s}^2} = 13 \text{ km}$$

A circle smaller than that will result
in pilot blackout.

(7) The equation of motion of the
particle is

$$m \frac{dv}{dt} = -mK (v^3 + a^2 v) \quad (1)$$

Integrating

$$\int \frac{dv}{v(\delta^2 + v^2)} = -k \int dt \quad (2)$$

$$\frac{1}{2\delta^2} \ln \left[\frac{v^2}{\delta^2 + v^2} \right] = -kt + C \quad (3)$$

Therefore, we have

$$\frac{v^2}{\delta^2 + v^2} = C' e^{-At} \quad (4)$$

where $A = 2\delta^2 k$ and C' is a new constant. We can evaluate C' by using the initial condition, $v = v_0$ at $t = 0$.

$$C' = \frac{v_0^2}{\delta^2 + v_0^2} \quad (5)$$

Substituting (5) into (4) and rearranging,

we have

$$v = \left[\frac{\delta^2 C' e^{-At}}{1 - C' e^{-At}} \right]^{1/2} = \frac{dx}{dt} \quad (6)$$

Now, in order to integrate (6), we introduce $u \equiv e^{-At}$ so that $du = -A u dt$

Then

$$x = \int \left[\frac{\partial^2 c' e^{-At}}{1 - c' e^{-At}} \right]^{1/2} dt = \frac{2}{A} \int \left[\frac{c' u}{1 - c' u} \right]^{1/2} \frac{du}{u}$$
$$= -\frac{\partial \sqrt{c'}}{A} \int \frac{du}{\sqrt{-c' u^2 + u}} \quad (7)$$

$$\text{Then } x = \frac{\partial}{A} \sin^{-1} (1 - 2c' u) + C'' \quad (8)$$

Again the constant C'' can be evaluated by putting $x=0$ at $t=0$; i.e., $x=0$

at $u=1$

$$C'' = -\frac{\partial}{A} \sin^{-1} (1 - 2c') \quad (9)$$

Therefore, we have

$$x = \frac{\partial}{A} \left[\sin^{-1} \left(-2c' e^{-At} + 1 \right) - \sin^{-1} (-2c' + 1) \right]$$

Using (4) and (5) we can write

$$x = \frac{1}{2k\alpha} \left[\sin^{-1} \left[\frac{-v^2 + \alpha^2}{v^2 + \alpha^2} \right] - \sin^{-1} \left[\frac{-v_0^2 + \alpha^2}{v_0^2 + \alpha^2} \right] \right] \quad (10)$$

From (6) we see that $v \rightarrow 0$ as $t \rightarrow \infty$
Therefore

$$\lim_{t \rightarrow \infty} \sin^{-1} \left[\frac{-v^2 + \alpha^2}{v^2 + \alpha^2} \right] = \sin^{-1}(1) = \pi/2 \quad (11)$$

Also for very large initial velocities

$$\lim_{v_0 \rightarrow \infty} \sin^{-1} \left[\frac{-v_0^2 + \alpha^2}{v_0^2 + \alpha^2} \right] = \sin^{-1}(-1) = -\pi/2 \quad (12)$$

Therefore using (10), (11), and (12)

$$x(t \rightarrow \infty) = \frac{\pi}{2k\alpha}$$

and the particle can never move a distance greater than $\pi/2k\alpha$ for any initial velocity.