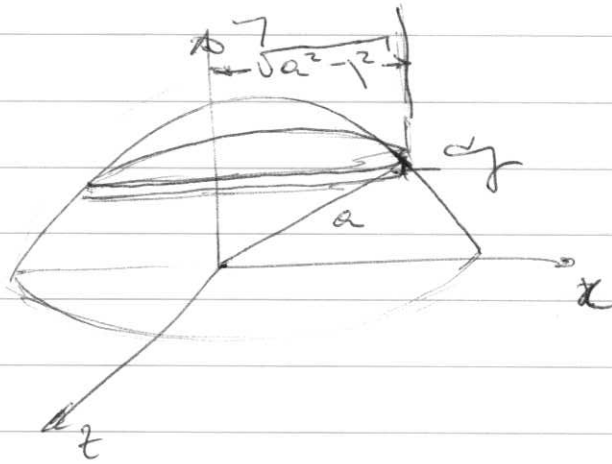


PROBLEMS # 11

(1)

Let the density be ρ , the hemisphere's mass be M , and the radius be a

$$\rho = \frac{M}{\frac{2}{3}\pi a^3}$$



The position coordinates of R_{CM} are (x, y, z) .

From symmetry $x=0$ $z=0$

$$x = \frac{1}{M} \int_{-a}^a x \, dm$$

$$y = \frac{1}{M} \int_{-a}^a z \, dm$$

$$Y = \frac{1}{M} \int_0^a y \, dm$$

$$dm = \rho \, dV = \rho \pi (a^2 - y^2) \, dy$$

$$Y = \frac{1}{M} \int_0^a \rho \pi (a^2 - y^2) y \, dy$$

$$Y = \frac{4\pi \rho a^4}{4M} = \frac{3a}{8}$$

$$R_{cm} = (0, 3a/8, 0)$$

(2) Put the steel in the $z > 0$ region with the base in the x, y plane.

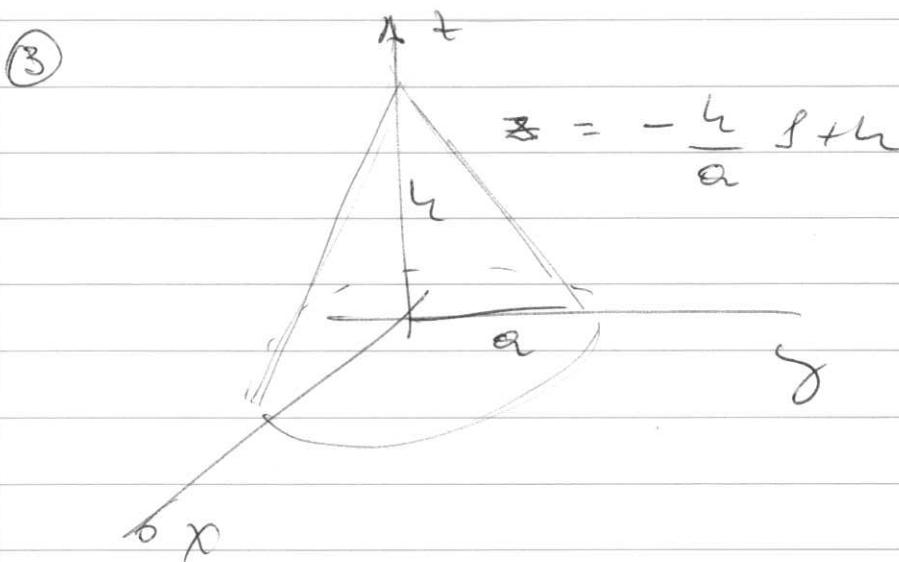
By symmetry $x_{cm} = y_{cm} = 0$

$$z = \frac{\int_0^{2\pi} \int_0^{\pi/2} \int_{r=0}^{r_2} \rho z r^2 \, dr \, d\theta \, dz}{\int_0^{2\pi} \int_0^{\pi/2} \int_{r=0}^{r_2} \rho r^2 \, dr \, d\theta \, dz}$$

Using $z = r \cos \theta$ and doing the integral

gives

$$z_{cm} = \frac{3}{8} \frac{(r_2^4 - r_1^4)}{(r_2^3 - r_1^3)}$$



By symmetry $x_{cm} = y_{cm} = 0$

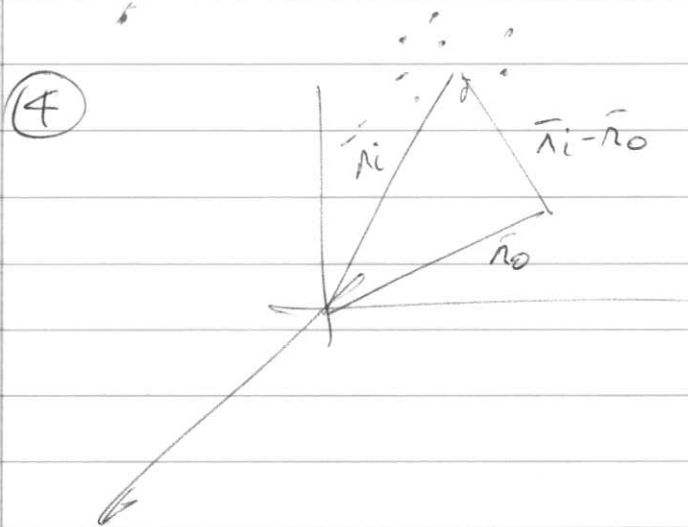
We use cylindrical coordinates

ρ = mass density

$$z_{cm} = \frac{\int_0^{2\pi} \int_0^a \int_0^{-\frac{h}{a}\rho+h} \rho z \rho d\rho d\phi dz}{\int_0^{2\pi} \int_0^a \int_0^{-\frac{h}{a}\rho+h} \rho d\rho d\phi dz} = \frac{h}{4}$$

The center-of-mass is on the axis of

the cone $\frac{3}{4} h$ from the vertex



r_i = position of the i^{th} particle

m_i = mass of the i^{th} particle

$M = \sum m_i$ = total mass

g = gravitational constant field

$$\tau = \sum \tau_i$$

$$= \sum (\vec{r}_i - \vec{r}_0) \times \vec{F}_i$$

$$= \sum (\vec{r}_i - \vec{r}_0) \times m_i g$$

$$= \sum r_i \times m_i g - \sum r_0 \times m_i g$$

$$= \sum m_i r_i \times g - (\sum m_i) r_0 \times g$$

$$= (\sum m_i r_i) \times g - M r_0 \times g$$

Now if the total torque is zero

$$\sum m_i r_i = M r_0$$

or

$$r_0 = \frac{1}{M} \sum m_i r_i$$

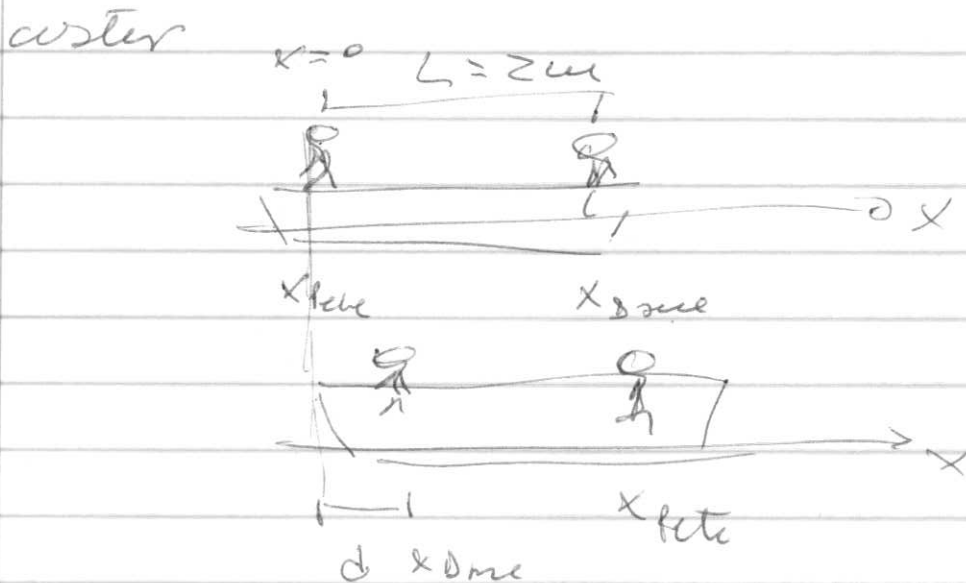
which is the definition of the center of mass.

So $\tau = 0$ about $r_0 = r_{CM}$, or

center of gravity = center of mass

(5) Flesh out $\Delta x_{cm} = \sum m_i x_i$ solve

before and after rate and mass change
 /less. the coordinate axis measures
 positions in the reference frame of the
 center



$$M x_{cm} = m_{pete} x_{pete} + m_{dave} x_{dave} + m_{sot} x_{sot}$$

$$\Delta x_{cm} = m_{pete} \Delta x_{pete} + m_{dave} \Delta x_{dave} + m_{sot} \Delta x_{sot}$$

$$\Delta \Delta x_{cm} = m_{pete} \Delta \Delta x_{pete} + m_{dave} \Delta \Delta x_{dave} + m_{sot} \Delta \Delta x_{sot}$$

$$0 = \mu_{\text{fete}} (d+L) + \mu_{\text{Dome}} (d+2L) + \mu_{\text{bst}} d$$

$$d = \frac{(\mu_{\text{Dome}} - \mu_{\text{fete}}) L}{\mu_{\text{Dome}} + \mu_{\text{fete}} + \mu_{\text{bst}}} \quad L = 0.31 \mu$$

$$\mu_{\text{Dome}} + \mu_{\text{fete}} + \mu_{\text{bst}}$$

(6) Let the velocity of the projectile of mass M be v . The three fragments have the following masses and velocities

$$m_1 = \frac{M}{2}, \quad v_1 = k_1 v \quad \text{Forward direction} \\ k_1 > 0$$

$$m_2 = \frac{M}{6}, \quad v_2 = -k_2 v \quad \text{opposite direction} \\ k_2 > 0$$

$$m_3 = \frac{M}{3}, \quad v_3 = 0 \quad \text{at rest}$$

The conservation of linear momentum and energy give

$$Mv = \frac{M}{2} k_1 v - \frac{M}{6} k_2 v \quad (1)$$

$$E + \frac{1}{2} M v^2 = \frac{1}{2} \frac{M}{2} (k_1 v)^2 + \frac{1}{2} \frac{M}{6} (k_2 v)^2 \quad (2)$$

From (1)

$$k_2 = 3k_1 - 6$$

which we can insert into (2)

$$5 \left(\frac{1}{2} M v^2 \right) + \frac{1}{2} M v^2 = \frac{M v^2}{4} k_1^2 + \frac{M v^2}{12} (3k_1 - 6)^2$$

which reduces to $k_1^2 - 3k_1 = 0$, giving the results $k_1 = 0$ and $k_1 = 3$. For $k_1 = 0$

the value of $k_2 = -6$, which is inconsistent

with the value of $k_2 > 0$. For $k_1 = 3$,

the value of $k_2 = 3$. The velocities become

$$v_1 = 3v$$

$$v_2 = -3v$$

$$v_3 = 0$$

(7) Since the initial kinetic energies of the two particles are equal we have

$$\frac{1}{2} m_1 u_1^2 = \frac{1}{2} m_2 u_2^2 = \frac{1}{2} \alpha^2 m_2 u_1^2 \quad (1)$$

$$\text{or } \frac{m_1}{m_2} = \alpha^2 \quad (2)$$

Now the kinetic energy of the system is conserved because the collision is elastic. Therefore

$$\frac{1}{2} m_1 u_1^2 + \frac{1}{2} m_2 u_2^2 = m_1 u_1^2 = \frac{1}{2} m_2 v_2^2 \quad (3)$$

Since $v_1 = 0$, momentum is also conserved

so we can write

$$m_1 u_1 + m_2 u_2 = (m_1 + \alpha m_2) u_1 = m_2 v_2 \quad (4)$$

Substituting the second equality in (4)

into (3) we find

$$m_1 u_1^2 = \frac{1}{2} m_2 \left[\frac{m_1 + \alpha m_2}{m_2} \right]^2 m_1^2 \quad (5)$$

or

$$m_1 = \frac{1}{2} m_2 \left[\frac{m_1 + \alpha}{m_2} \right]^2 \quad (4)$$

Using $m_1 m_2 = \alpha^2$ (4) becomes

$$2\alpha^2 = (\alpha^2 + \alpha)^2$$

solving for α we obtain

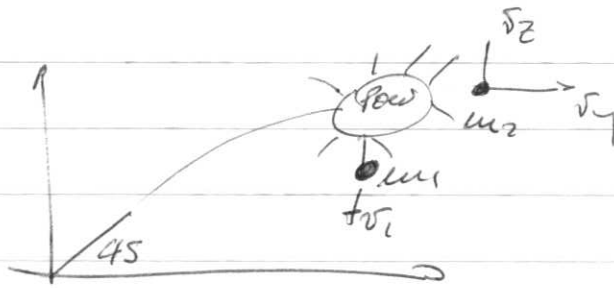
$$\alpha_1 = -1 \pm \sqrt{2}$$

$$\alpha_2 = 3 \mp 2\sqrt{2}$$

But given in

$$\frac{m_1}{m_2} = 3 \pm 2\sqrt{2}, \quad \frac{u_2}{u_1} = -(1 \pm \sqrt{2}) \text{ with } \begin{cases} +: \alpha < 0 \\ -: \alpha > 0 \end{cases}$$

(8)



Let the axis be as shown with the projectile in the $y-z$ plane. At the top just before the explosion, the velocity is in the y direction and has magnitude

$$v_{0y} = \frac{v_0}{\sqrt{2}}$$

$$v_{0y} = \frac{v_0}{\sqrt{2}} = \frac{\sqrt{\frac{2E_0}{m_1+m_2}}}{\sqrt{2}} = \sqrt{\frac{E_0}{m_1+m_2}}$$

where m_1 and m_2 are the masses of the fragments. The initial momentum is

$$P_i = (m_1 + m_2) \left[0, \sqrt{\frac{E_0}{m_1+m_2}}, 0 \right]$$

The final momentum is

$$P_f = P_1 + P_2$$

$$P_1 = m_1 (0, 0, v_1)$$

$$P_2 = m_2 (v_x, v_y, v_z)$$

The conservation of momentum

equations are

$$P_x: 0 = m_2 v_x \quad \text{or} \quad v_x = 0$$

$$P_y: \sqrt{E_0 (m_1 + m_2)} = m_2 v_y$$

$$\text{or} \quad v_y = \frac{1}{m_2} \sqrt{E_0 (m_1 + m_2)}$$

$$P_z: 0 = m_1 v_1 + m_2 v_z$$

$$v_1 = -\frac{m_2}{m_1} v_z$$

The energy equation is

$$\frac{1}{2} (m_1 + m_2) \frac{E_0}{(m_1 + m_2)} + E_0 = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 (v_y^2 + v_z^2)$$

or

$$3E_0 = m_1 v_1^2 + m_2 (v_y^2 + v_z^2)$$

Substituting for v_y and v_1 gives

$$v_z = \sqrt{\frac{E_0 m_1 (2m_2 - m_1)}{m_2^2 (m_1 + m_2)}}$$

$$v_1 = -\frac{m_2}{m_1} v_z \text{ gives}$$

$$v_1 = -\sqrt{\frac{E_0 (2m_2 - m_1)}{m_1 (m_1 + m_2)}}$$

So m_1 travels straight down with speed $|v_1|$

v_2 travels in the $y-z$ plane

$$v_2 = (v_y^2 + v_z^2)^{1/2} = \sqrt{\frac{E_0 (4m_1 + m_2)}{m_2 (m_1 + m_2)}}$$

$$\theta = \arctan \frac{v_z}{v_y} = \arctan \frac{\sqrt{m_1 (2m_2 - m_1)}}{(m_1 + m_2)}$$

(9) Conservation of momentum requires v_f to be in the same direction as u_1 (component of $v_f \perp$ to u_1 must be zero).

$$p_i = m_1 u_1$$

$$p_f = (m_1 + m_2) v_f$$

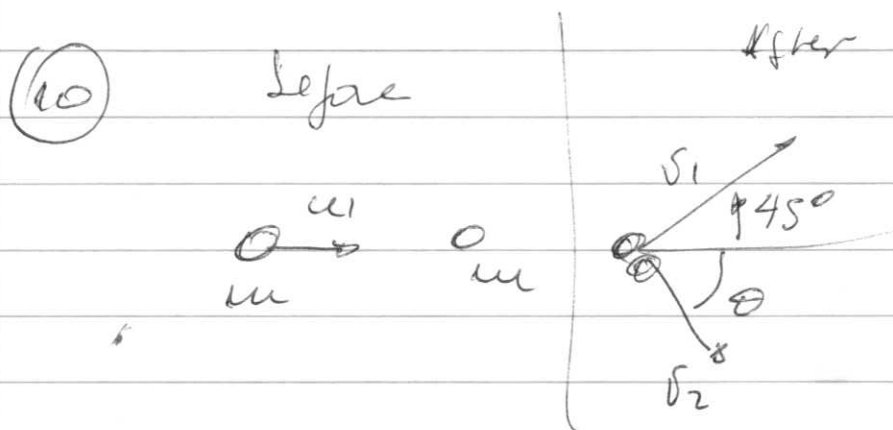
$$p_i = p_f \Rightarrow v_f = \frac{m_1}{m_1 + m_2} u_1$$

The fraction of the original kinetic energy lost is

$$\frac{K_i - K_f}{K_i} = \frac{\frac{1}{2} m_1 u_1^2 - \frac{1}{2} (m_1 + m_2) \frac{m_1^2 u_1^2}{(m_1 + m_2)^2}}{\frac{1}{2} m_1 u_1^2}$$

$$= \frac{m_1 - m_1^2 / (m_1 + m_2)}{m_1}$$

$$= \frac{m_2}{m_1 + m_2}$$



Conservation of P_x :

$$mu u_1 = mu v_1 \cos 45^\circ + mu v_2 \cos \theta \quad (1)$$

Conservation of P_y :

$$0 = mu v_1 \sin 45^\circ - mu v_2 \sin \theta \quad (2)$$

Conservation of energy (elastic collision)

$$\frac{1}{2} mu u_1^2 = \frac{1}{2} mu v_1^2 + \frac{1}{2} mu v_2^2 \quad (3)$$

Solve (1) for $\cos \theta$

$$\cos \theta = \frac{u_1 - v_1 / \sqrt{2}}{v_2}$$

Solve (2) for $\sin \theta$

$$\sin \theta = \frac{v_1}{\sqrt{2} v_2}$$

Substitute into $\cos^2 \theta + \sin^2 \theta = 1$, $\sin \theta = u_1 / \sqrt{2}$

and the result is

$$u_1^2 = \sqrt{2}^2 - \sqrt{2}^2 + \sqrt{2} u_1 \sqrt{2}$$

Combining that with (3) gives

$$2\sqrt{2}^2 = \sqrt{2} u_1 \sqrt{2}$$

we know $\sqrt{2} \neq 0 \Rightarrow$

$$\sqrt{2} = u_1 / \sqrt{2}$$

Substitute into (3) and the result is

$$\sqrt{2} = u_1 / \sqrt{2}$$

Since $\sqrt{2} = \sqrt{2}$, (2) implies that

$$\theta = 45^\circ$$

(11)

Using $y = v_0 t - \frac{1}{2} g t^2$ and $v = v_0 - g t$,
we can get the velocities before and after

the collision

Before $u_1 = -g t_1$ where $h_1 = \frac{1}{2} g t_1^2$

$$\text{So } u_1 = -g \sqrt{\frac{2h_1}{g}} = -\sqrt{2gh_1}$$

After $0 = v_0 - g t_2$ or $t_2 = v_0/g$

$$h_2 = v_0 t_2 - \frac{1}{2} g t_2^2$$

$$= \frac{v_0^2}{g} - \frac{1}{2} \frac{v_0^2}{g} \quad \text{or } v_0 = \sqrt{2gh_2}$$

$$\text{So } v_1 = \sqrt{2gh_2}$$

Thus

$$e = \frac{|v_2 - v_1|}{|u_2 - u_1|} = \frac{\sqrt{2gh_2}}{\sqrt{2gh_1}}$$

$$e = \sqrt{\frac{h_2}{h_1}}$$

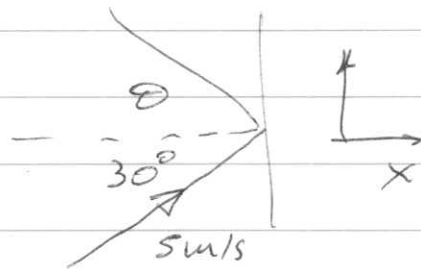
$$T_{\text{lost}} = T_i - T_f$$

$$\text{fraction lost} = \frac{T_i - T_f}{T_i}$$

$$= \frac{u_1^2 - v_1^2}{u_1^2} = \frac{h_1 - h_2}{h_1} = 1 - \frac{h_2}{h_1}$$

$$\frac{T_i - T_f}{T_i} = 1 - e^2$$

(12)



The velocity component in the y direction is unchanged

$$v_y = u_y = 5 \frac{\text{m}}{\text{s}} \sin 30^\circ = 2.5 \text{ m/s}$$

For the x component we have

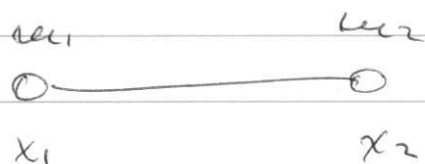
$$0.8 = \frac{|\delta_x|}{|u_x|} = \frac{\delta_x}{\left(\frac{5 \text{ m}}{\text{s}}\right) \cos 30^\circ} = \frac{\delta_x}{\frac{\sqrt{3}}{2} 5 \text{ m/s}}$$

$$|\delta_x| = 2\sqrt{3} \text{ m/s}$$

$$v_f = \frac{1}{2} \sqrt{73} \approx 4.3 \text{ m/s}$$

$$\theta = \text{ARCTAN} \frac{2.5}{2\sqrt{3}} \approx 36^\circ$$

13



$$r = |\vec{x}_2 - \vec{x}_1|$$

when two particles are initially at rest separated by a distance r_0 , the system has the total energy

$$E_0 = -G \frac{m_1 m_2}{r_0} \quad (1)$$

The coordinates of the particles, x_1 and x_2 , are measured from the position of the center of mass. At any time the total energy is

$$E = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_2^2 - \frac{G m_1 m_2}{r} \quad (2)$$

and the linear momentum, at any time, is

$$p = m_1 \dot{x}_1 + m_2 \dot{x}_2 = 0 \quad (3)$$

From the conservation of energy we have

$$E = E_0, \text{ or}$$

$$- \frac{G m_1 m_2}{r_0} = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_2^2 - \frac{G m_1 m_2}{r} \quad (4)$$

Using (3) in (4), we find

$$\dot{x}_1 = v_1 = u_2 \sqrt{\frac{2G}{M} \left[\frac{1}{r} - \frac{1}{r_0} \right]}$$

$$\dot{x}_2 = v_2 = -u_1 \sqrt{\frac{2G}{M} \left[\frac{1}{r} - \frac{1}{r_0} \right]}$$

(14)

$$(a) \quad v = v_0 + u \ln \frac{u_0}{u}$$

Assuming $v_0 = 0$, we have

$$v = \left[100 \frac{\text{m}}{\text{s}} \right] \ln \frac{100}{98}$$

$$v = 2.02 \text{ m/s}; \text{ so, he runs out}$$

of gas

(b) Relative to Subject original frame of reference we have

Before throwing back

$$\boxed{98 \text{ kg}} \rightarrow 2.02 \text{ m/s}$$

After throwing back we want Subject velocity to be slightly greater than $3 \frac{\text{m}}{\text{s}}$

(so he will catch up the orbiter)

$$v \leftarrow \boxed{8 \text{ kg}} \quad \boxed{90 \text{ kg}} \rightarrow 3 \frac{\text{m}}{\text{s}}$$

Conservation of momentum gives

$$98 \text{ kg} (2.02 \text{ m/s}) = (90 \text{ kg})(3 \text{ m/s}) - (8 \text{ kg})v$$

$$v = 9 \text{ m/s}$$

(This velocity is relative to Subject reference frame, i.e., before he fires his permissoid tank.) Since Subject

is travelling towards the orbiter at $2.02 \frac{\text{m}}{\text{s}}$,

he must throw the tank at $v = 9 \text{ m/s} + 2.02 \text{ m/s}$

$$v = 11 \text{ m/s}$$