

# Group Problems #27 - Solutions

Wednesday, November 2

## Problem 1 *Gaussian solution to the harmonic oscillator*

As we'll see, the ground-state of the quantum harmonic oscillator has a Gaussian wavefunction of the form  $\psi(x) = Ae^{-ax^2}$ .

- (a) What are the units (dimensions) of the constant  $a$ ?

The argument of the exponential function must be dimensionless (otherwise, what would be the dimension of the exponential function itself?). Thus,  $a$  must have units of  $1/\text{length}^2$ :  $a \equiv L^{-2}$ .

- (b) Use the kinetic energy operator  $\hat{K}$  to find an expression for the constant  $a$  in terms of the classical turning point,  $x_T$ , where a classical oscillator (e.g., mass on a spring) would change direction.

The kinetic energy operator is defined as:

$$\hat{K} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}, \quad (1)$$

where  $\hat{K}$  acting on the wave function gives the *value* of the kinetic energy  $K(x)$  at a particular position  $x$  multiplied by the wave function:  $\hat{K}\psi(x) = K(x)\psi(x)$ . Applying this to the wave function above gives:

$$\hat{K}\psi(x) = K(x)\psi(x) \implies -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} [Ae^{-ax^2}] = \frac{\hbar^2 a}{m} Ae^{-ax^2} [1 - 2ax^2] \quad (2)$$

$$\implies K(x) = \frac{\hbar^2 a}{m} [1 - 2ax^2]. \quad (3)$$

At the classical turning point,  $x = x_T$ , the kinetic energy is zero, so

$$K(x = x_T) = 0 \implies 1 - 2ax_T^2 = 0 \quad (4)$$

$$\implies a = \frac{1}{2x_T^2}, \quad (5)$$

and we see that  $a$  has units of  $1/\text{length}^2$  as required. Note that  $E = U(x)$  at  $x = x_T$ , and there is an inflection point in  $\psi(x)$ .

- (c) Consider a particle in the ground state of the harmonic oscillator potential. If you were to make a measurement of its kinetic energy  $K$  at a particular point in space, then you must get a number. If you repeat this measurement over and over again, will you get the same number? Why or why not? (*Hint*: think about the commutation relationship between the position operator  $\hat{x}$  and the kinetic energy operator  $\hat{K}$ .)

Using the hint, we should calculate the commutator of  $\hat{K}$  and  $\hat{x}$ :  $[\hat{K}, \hat{x}] = \hat{K}\hat{x} - \hat{x}\hat{K}$ , where:

$$\hat{K} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \quad (6)$$

$$\hat{x} = x. \quad (7)$$

Applying this to the wave function, we have:

$$\hat{K}\hat{x}\psi(x) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} [x A e^{-ax^2}] \quad (8)$$

$$(9)$$

$$= -A \frac{\hbar^2}{2m} \frac{\partial}{\partial x} [e^{-ax^2} (1 - 2ax^2)] \quad (10)$$

$$(11)$$

$$= A e^{-ax^2} \frac{\hbar^2}{2m} 2ax (3 - 2ax^2), \quad (12)$$

and

$$\hat{x}\hat{K}\psi(x) = -\frac{\hbar^2}{2m} x \frac{\partial^2}{\partial x^2} [A e^{-ax^2}] \quad (13)$$

$$(14)$$

$$= \frac{\hbar^2}{2m} Ax \frac{\partial}{\partial x} [2axe^{-ax^2}] \quad (15)$$

$$(16)$$

$$= A e^{-ax^2} \frac{\hbar^2}{2m} 2ax (1 - 2ax^2). \quad (17)$$

Putting this together gives:

$$[\hat{K}, \hat{x}]\psi(x) = [\hat{K}\hat{x} - \hat{x}\hat{K}]Ae^{-ax^2} \quad (18)$$

$$(19)$$

$$= A e^{-ax^2} \frac{\hbar^2}{2m} 2ax [(3 - 2ax^2) - (1 - 2ax^2)] \quad (20)$$

$$(21)$$

$$= A e^{-ax^2} \frac{\hbar^2}{2m} 2ax \cdot 2 \quad (22)$$

$$(23)$$

$$\implies [\hat{K}, \hat{x}] = \frac{2a\hbar^2}{m} x \neq 0. \quad (24)$$

Since this is not equal to zero, the  $\hat{K}$  and  $\hat{x}$  do not commute, and we cannot simultaneously measure the particle's kinetic energy and position simultaneously. So if we constrain our measurement to a particular value of position ( $x$ ), then we will measure a spread in kinetic energy values when we repeat the measurement many times on a similarly prepared system.

- (d) What is the expectation value for the momentum of a particle in the ground state of the harmonic oscillator? (*Hint:* This question involves doing an apparent integral, but you can use symmetry arguments to avoid actually computing the integral.)

The expectation value of the physical quantity (observable) associated with operator  $\hat{O}$  is given by:

$$\langle O \rangle = \int_{-\infty}^{\infty} \psi^*(x) \hat{O} \psi(x) dx. \quad (25)$$

In this case,  $\hat{O} = \hat{p} = -i\hbar \frac{\partial}{\partial x}$ . So finally,

$$\langle p \rangle = -i\hbar A^2 \int_{-\infty}^{\infty} e^{-ax^2} \frac{d}{dx} (e^{-ax^2}) dx \quad (26)$$

$$(27)$$

$$= 2i\hbar A^2 a \int_{-\infty}^{\infty} x e^{-2ax^2} dx = 0, \quad (28)$$

since the argument is an odd function of  $x$  integrated over symmetric bounds. Thus, the expectation value of the momentum is zero: this doesn't mean that a measurement of the momentum will equal zero. It means that repeated measurements of the momentum on similarly prepared systems will yield a mean value of zero.