

1. Show that the charge  $Q = \int d^3x j^0$  must be a conserved quantity because of  $U(1)$  phase invariance.

2. The Lagrangian for three interacting real fields  $\phi_1, \phi_2, \phi_3$  is

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \phi_i)^2 - \frac{1}{2}\mu^2 \phi_i^2 - \frac{1}{4}\lambda(\phi_i^2)^2, \quad (1)$$

where  $\mu^2 < 0$ ,  $\lambda > 0$ , and a summation of  $\phi_i^2$  over  $i$  is implied. Show that it describes a massive field of mass  $\sqrt{-2\mu^2}$  and two massless Goldstone bosons.

3. Rather than (2.4.60), take instead  $\phi$  to be an  $SU(2)$  triplet of real scalar fields. For  $\mu^2 < 0$  and  $\lambda > 0$ , show that in this case two gauge bosons acquire mass but that the third remains massless. [Hint: Verify, and use,  $(T_k)_{ij} = -i\epsilon_{ijk}$  for the triplet representation of  $SU(2)$ .]

4. Generic extensions of the standard model predict the existence of extra  $U(1)$  gauge symmetries beyond hypercharge. However, in order to avoid long range forces other than gravity and Coulomb forces, only  $U(1)$  of hypercharge must survive as a massless gauge boson and all other extra Abelian gauge bosons must grow a mass. To understand the basis of the Stueckelberg mechanism giving mass to the  $U(1)$ 's, consider the following Lagrangian

$$\mathcal{L} = -\frac{1}{12}H^{\mu\nu\rho}H_{\mu\nu\rho} - \frac{1}{4g^2}F^{\mu\nu}F_{\mu\nu} + \frac{c}{4}\epsilon^{\mu\nu\rho\sigma}B_{\mu\nu}F_{\rho\sigma}$$

coupling an Abelian gauge field  $A_\mu$  to an antisymmetric tensor  $B_{\mu\nu}$ , where

$$H_{\mu\nu\rho} = \partial_\mu B_{\nu\rho} + \partial_\rho B_{\mu\nu} + \partial_\nu B_{\rho\mu}, \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu,$$

and  $g, c$  are arbitrary constants. This corresponds to the kinetic term for the fields  $B_{\mu\nu}$  and  $A_\mu$  together with the  $B \wedge F$  term.

(i) Imposing the constraint  $H = dB$  by the standard introduction of a Lagrange multiplier field  $\eta$ , rewrite the Lagrangian in terms of the (arbitrary) field  $H_{\mu\nu\rho}$ .

(ii) Show that the Lagrangian is dual to

$$\mathcal{L} = -\frac{1}{4g^2}F^{\mu\nu}F_{\mu\nu} - \frac{c^2}{2}(A_\sigma + \partial_\sigma \eta)^2,$$

which is just a Stueckelberg mass term for the gauge field  $A_\mu$  after “eating” the scalar  $\eta$  to acquire a mass  $m^2 = g^2 c^2$ .

5. Show that

$$\frac{1}{2}\tau^3 L_L = T^3 L_L,$$

where

$$L_L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}, \quad \tau^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix},$$

and

$$T^3 = \frac{1}{2} \int \psi_L^\dagger \tau^3 \psi_L d^3x.$$