

Astronomy, Astrophysics, and Cosmology

Luis A. Anchordoqui

Department of Physics and Astronomy
Lehman College, City University of New York

Lesson X
April 19, 2016

[arXiv:0706.1988](https://arxiv.org/abs/0706.1988)

At the beginning... there was the hand of God



Table of Contents

- 1 The Early Universe
 - The first millisecond

- History of universe from 10^{-10} seconds to today
is based on observational facts
- Fundamental laws of high energy physics are well-established
up to energies reached by LHC
- Before 10^{-10} seconds \Rightarrow energy of universe exceeds 13 TeV
and we lose comfort of direct experimental guidance
- Physics of that era is as speculative as it is fascinating
- Today \Rightarrow we will go back to the earliest of times
as close as possible to the *big bang*
and follow evolution of Universe

- As $a \rightarrow 0$ \Rightarrow temperature increases without limit $T \rightarrow \infty$
but there comes point @ which extrapolation of classical physics
breaks down
- Realm of quantum black holes \Rightarrow thermal energy of particles
is such that their de Broglie wavelength
is smaller than their Schwarzschild radius
- Equating h/mc to $2Gm/c^2$ yields system of Planck units

$$\begin{aligned}
 M_{\text{Pl}} &\equiv \sqrt{\frac{\hbar c}{G}} \simeq 10^{19} \text{ GeV} \\
 \ell_{\text{Pl}} &\equiv \sqrt{\frac{\hbar G}{c^3}} \simeq 10^{-35} \text{ m} \\
 t_{\text{Pl}} &\equiv \sqrt{\frac{\hbar G}{c^5}} \simeq 10^{-43} \text{ s}
 \end{aligned}
 \tag{1}$$

- t_{Pl} \Rightarrow sets origin of time for *classical big bang* era
- It is inaccurate to extend solution of Friedmann equation to $a = 0$
and conclude universe began in singularity of infinite density


- @ $t \sim 10^{-43}$ s \Rightarrow *phase transition* is thought to have occurred during which gravitational force *condensed out* as separate force
- Symmetry of four forces was broken but the strong, weak, and electromagnetic forces were still unified and there were no distinctions between quarks and leptons
- This is unimaginably short time
and predictions can be only speculative
- Temperature would have been about 10^{32} K \Rightarrow corresponding to *particles* moving about every which way with average energy

$$kT \approx \frac{1.4 \times 10^{-23} \text{ J/K } 10^{32} \text{ K}}{1.6 \times 10^{-10} \text{ J/GeV}} \approx 10^{19} \text{ GeV} \quad (2)$$

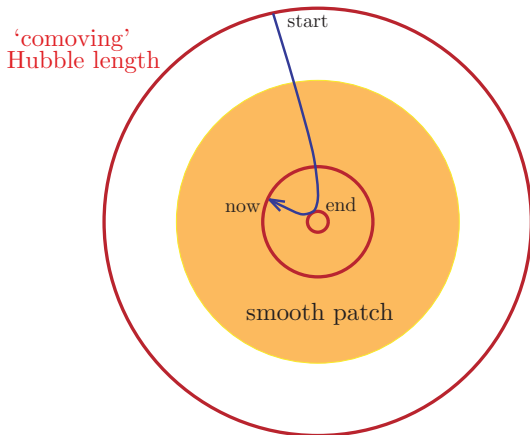
- Very shortly thereafter as temperature had dropped to $\sim 10^{28}$ K there was another phase transition
where strong force condensed out
- Now universe was filled with *soup* of quarks and leptons
- About this time universe underwent exponential expansion increasing in size by a factor of $\gtrsim 10^{26}$
in a tiny fraction of a second perhaps $\sim 10^{-34}$ s
- Favored Λ CDM model implicitly includes inflationary hypothesis where scale factor expands exponentially $a(t) \propto e^{Ht}$
- If interval of exponential expansion satisfies $\Delta t \gtrsim 60/H$
small casually connected region can grow sufficiently
to accommodate observed homogeneity and isotropy

- Important to emphasize subtle distinction between comoving horizon ϱ_h and comoving Hubble radius $c/(aH)$
- Express comoving horizon as integral of comoving Hubble radius

$$\varrho_h \equiv c \int_0^t \frac{dt'}{a(t')} = c \int_0^a \frac{da}{Ha^2} = c \int_0^a \frac{1}{aH} d \ln a \quad (3)$$

- If particles are separated by distances greater than ϱ_h they never could have communicated with one another
- If they are separated by distances greater than $c/(aH)$ they cannot talk to each other now
- This distinction is crucial for solution to horizon problem
- It is possible that ϱ_h is much larger than $c/(aH)$ now so that particles cannot communicate today but were in causal contact early on
- From (3) we see that this might happen if comoving Hubble radius in early universe \gg that it is now so that ϱ_h got most of its contribution from early times
- Hence  we require phase of decreasing Hubble radius

- Evolution of comoving Hubble radius in inflationary universe



- Comoving Hubble sphere shrinks during inflation
and expands after inflation
- Inflation \Rightarrow mechanism to *zoom-in* on smooth sub-horizon patch

Conditions for inflation

- shrinking Hubble sphere defined by

$$\frac{d}{dt} \left(\frac{1}{aH} \right) < 0 \quad (4)$$

- From \Rightarrow

$$\frac{d}{dt} \left(\frac{1}{aH} \right) = -\frac{\ddot{a}}{(aH)^2} \quad (5)$$

shrinking comoving Hubble radius

implies accelerated expansion $\ddot{a} > 0$

- This explains why
inflation is often defined as period of accelerated expansion

More conditions for inflation

- Second time derivative of scale factor can be related to first time derivative of Hubble parameter

$$\frac{\ddot{a}}{a} = H^2(1 - \epsilon) \Rightarrow \epsilon \equiv -\frac{\dot{H}}{H^2} \quad (6)$$

- Acceleration therefore corresponds to $\epsilon < 1$
- During inflation
 - H is approximately constant during inflation
 - a grows exponentially
 - this implies $c/(aH)$ decreases \Rightarrow just as advertised
- Consulting *acceleration equation*

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3c^2}(\rho + 3P). \quad (7)$$

we infer that $\ddot{a} > 0$ requires a negative pressure $\Rightarrow P < -\rho/3$

- After very brief inflationary period universe would have settled back into its more regular expansion
- For $10^{-34} \text{ s} < t < 10^5 \text{ yr} \Leftrightarrow 10^3 \text{ K} < T < 10^{27} \text{ K}$ universe is thought to have been dominated by radiation
- We have seen equation of state is given by $w = 1/3$
- Neglect contributions to H from Λ
(this is always a good approximation for small enough a)
then $\Leftrightarrow a \sim t^{1/2}$ and $\rho_{\text{rad}} \sim a^{-4}$
- Expansion rate as a function of temperature in plasma

$$\begin{aligned}
 H &= \left(\frac{8\pi G \rho_{\text{rad}}}{3} \right)^{1/2} = \left(\frac{8\pi^3}{90} g_{\rho}(T) \right)^{1/2} T^2 / M_{\text{Pl}} \\
 &\sim 1.66 \sqrt{g_{\rho}(T)} T^2 / M_{\text{Pl}} \quad (8)
 \end{aligned}$$

(we have adopted natural units $\hbar = c = k = 1$)


- Neglecting T -dependence of g_ρ
(i.e. away from mass thresholds and phase transitions)
integration of (8) yields

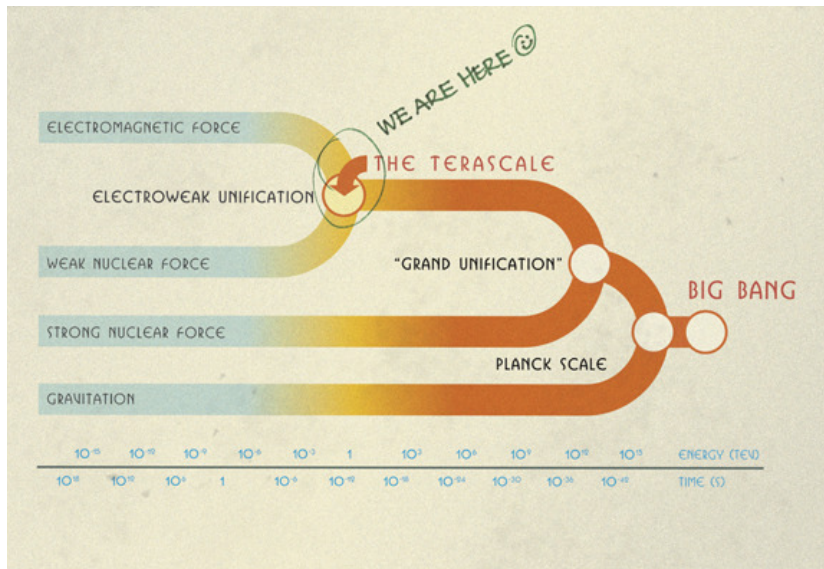
$$a(t) = \left(\frac{t}{t_0}\right)^{1/2} \quad \text{and} \quad \rho(t) = \frac{\rho_0}{a^4} = \frac{\rho_0 t_0^2}{t^2} \quad (9)$$

- (8) leads to useful commonly used approximation

$$t \simeq \left(\frac{3M_{\text{Pl}}^2}{32\pi\rho_{\text{rad}}}\right)^{1/2} \simeq 2.42 \frac{1}{\sqrt{g_\rho}} \left(\frac{T}{\text{MeV}}\right)^{-2} \text{ s} \quad (10)$$

Electroweak symmetry breaking

- At about 10^{-10} s  Higgs field spontaneously acquires VEV which breaks electroweak gauge symmetry
- Weak force and electromagnetic force manifest with different ranges
- Through Higgs mechanism quarks and charged leptons become massive
- Fundamental interactions have then taken their present forms but the temperature of the universe ($T \sim 1$ TeV) is still too high to allow quarks to bind together to form hadrons

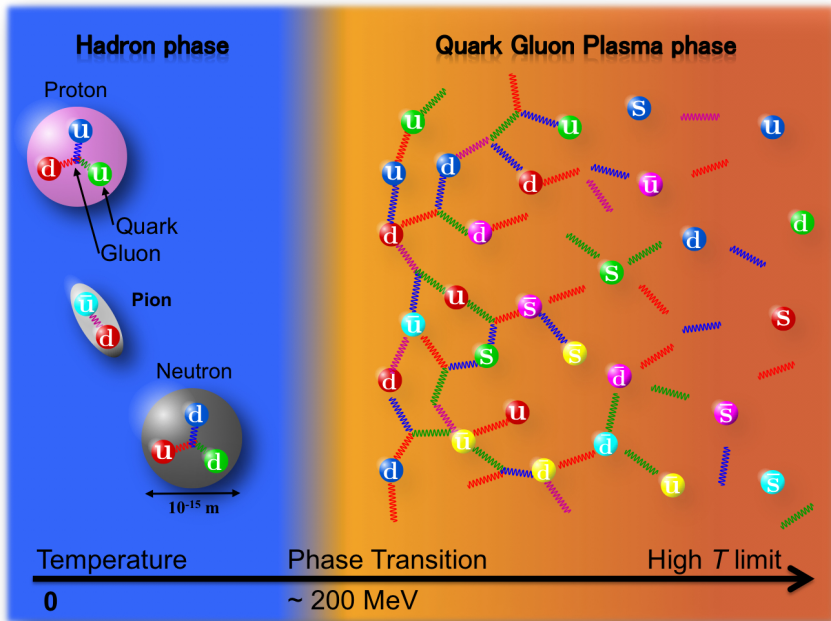


- @ $t \sim 10^{-6}$ s $\Rightarrow T \sim 1$ GeV
quarks began to *condense* into mesons and baryons
- If relativistic particles are present that have decoupled from γ 's it is necessary to distinguish between two kinds of r.d.o.f.:
 - those associated total energy density g_ρ
 - those associated with total entropy density g_s
- @ energies above deconfinement transition
quarks and gluons are relevant fields for QCD sector
effective number of interacting (thermally coupled) r.d.o.f. is

$$g_s(T) = 61.75$$

- As universe cools below confinement scale $\Rightarrow \Lambda_{\text{QCD}} \sim 200$ MeV
SM plasma transitions to regime
where mesons and baryons are pertinent degrees of freedom
- Precisely \Rightarrow relevant hadrons present in this energy regime
are pions and charged kaons

$$g_s(T) = 19.25$$



- Significant reduction in r.d.o.f. is from rapid annihilation or decay of massive hadrons which may have formed during transition
- Quark-hadron crossover transition \Rightarrow has associated large redistribution of entropy into remaining r.d.o.f.
- To connect temperature to effective number of r.d.o.f. use high statistics lattice simulations of QCD plasma in hot phase especially behavior of entropy during changeover
- Concretely \Rightarrow effective number of interacting r.d.o.f. in plasma

$$g_s(T) \simeq r(T) \left(g_B + \frac{7}{8} g_F \right) \quad (11)$$

◇ $r = 1$ for leptons

◇ $r = 22$ for photon contributions

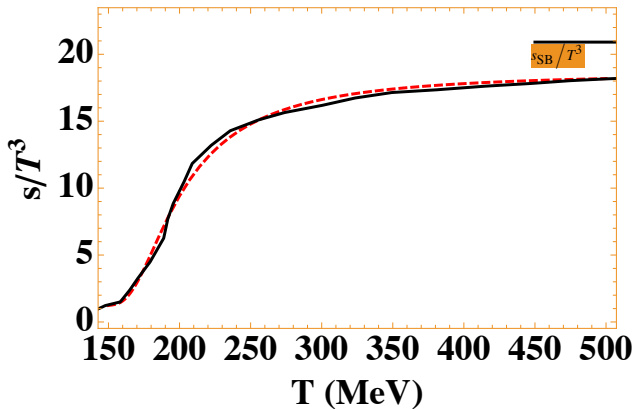
◇ $r = s(T)/s_{SB}$ for quark-gluon plasma

- $s(T)$ \Rightarrow actual entropy
- (s_{SB}) \Rightarrow ideal Stefan-Boltzmann entropy

- ✧ For $150 \text{ MeV} < T < 500 \text{ MeV}$
 entropy rise during confinement-deconfinement changeover

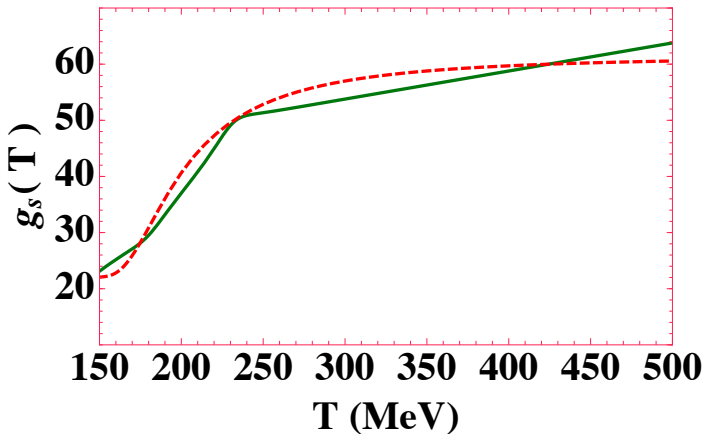
$$\frac{s}{T^3} \simeq \frac{42.82}{\sqrt{392\pi}} e^{-C_1} + 18.62 \frac{C_2^2}{[e^{C_2} - 1]^2} e^{C_2} \quad (12)$$

$$C_1 = (T_{\text{MeV}} - 151)^2 / 392 \text{ and } C_2 = 195.1 / (T_{\text{MeV}} - 134)$$



✧ For same energy range

$$g_s(T) \simeq 47.5 r(T) + 19.25 \quad (13)$$




- Entropy density dominated by contribution of relativistic particles
- To very good approximation

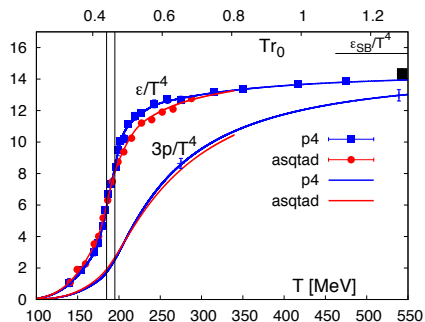
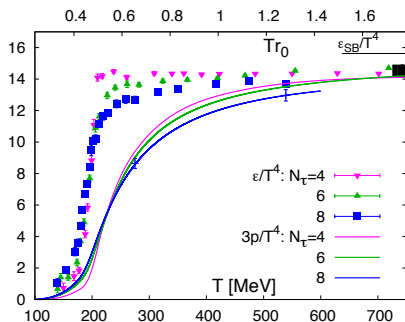
$$s = \frac{2\pi^2}{45} g_s(T) T^3 \quad (14)$$

- Conservation of $S = sV$ leads to

$$\frac{d}{dt}(sa^3) = 0 \Rightarrow g_s(T)T^3a^3 = \text{constant as universe expands} \quad (15)$$

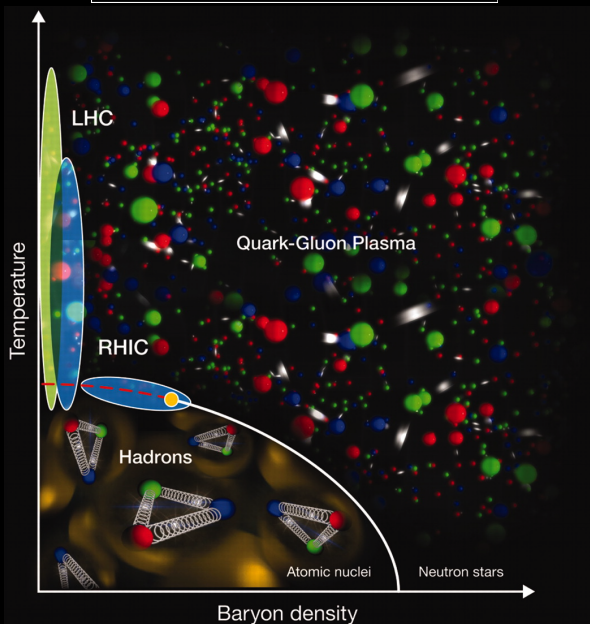
- As one would expect  non-evolving system would stay @ constant entropy density in comoving coordinates even though s is decreasing due to expansion of universe

- ✧ Since quark-gluon energy density in plasma has similar T dependence to that of the entropy



we'll simplify discussion by taking $g = g_\rho = g_s$

Sensitivity of particle colliders



Baryogenesis

- Manned and unmanned exploration of solar system tell us that it is made up of same stuff as Earth: *baryons*
- Observational evidence from radio-astronomy and cosmic rays indicate Milky Way and distant galaxies are made of baryons
- Conclusion: baryon number of the observable universe $\Rightarrow B > 0$
- Requirement: early $q\bar{q}$ plasma contained tiny surplus of quarks
- After all anti-matter annihilated with matter
only small surplus of matter remained

$$\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma} = 5 \times 10^{-10} \frac{\text{excess baryons}}{\text{photons}} \quad (16)$$

- Tiny surplus \Rightarrow explained by interactions in early universe that were not completely symmetric with respect to an exchange of matter-antimatter