

PHYSICS 307



MATHEMATICAL PHYSICS

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ORDINARY DIFFERENTIAL EQUATIONS V

3.1 Setting the Stage ✓

3.2 Initial Value Problem ✓

Picard's existence and uniqueness theorem

Systems of first-order linear differential equations

Green matrix as a generalized function

3.3 Boundary Value Problem ✓

Self-adjointness of Sturm-Liouville operator

Green function of Sturm-Liouville operator

Series solutions to homogeneous linear equations

3.4 Fourier Analysis

Fourier series

Fourier transform

Fourier Analysis

Consider Sturm-Liouville eigenvalue problem $L = -d^2/dx^2$, $\rho = 1$

$$f(-\pi) = 0 \quad \text{and} \quad f(\pi) = 0 \quad (3.4.316.)$$

eigenfunctions are $\leftarrow b_n \sin(nx)$ with $n \in \mathbb{Z}$

Sturm-Liouville theorem states

if $f(x)$ is of class $C^2[-\pi, \pi]$ and satisfies (3.4.316.)

it can be expanded in a convergent series

$$f(x) = \sum_{n=1}^{\infty} b_n \sin(nx)$$

if we replace boundary conditions $\leftarrow f'(-\pi) = 0$ and $f'(\pi) = 0$

We are now speaking of quite a different Sturm-Liouville system

in this space eigenfunctions are $\leftarrow a_n \cos(nx)$ with $n \in \mathbb{Z}$

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(nx)$$

Fourier series

Central idea of today's class

investigate to what extent it is possible

to expand $f(x) \notin C^2[-\pi, \pi]$ in an infinite series

$$f(x) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} [a_n \cos(nx) + b_n \sin(nx)] \quad (3.4.319.)$$

1st order of business is to determine coefficients a_n and b_n

only then will we deal with convergence issues


Definition 3.4.1.

Assuming expansion (3.4.319.) holds we can determine a_n and b_n using orthogonality (and norm) of $\sin(nx)$ and $\cos(nx)$ integrating on both sides of (3.4.319.) over $[-\pi, \pi]$ we obtain

$$\int_{-\pi}^{\pi} f(x) dx = \pi a_0 \quad (3.4.320.)$$

note that
$$\frac{a_0}{2} = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx \quad (3.4.321.)$$

is average value of $f(x)$ in the interval $[-\pi, \pi]$

To calculate a_n with $n \neq 0$ 

multiply both sides of (3.4.319) by $\cos(kx)$, $k \neq 0$ and integrate

$$a_k = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos(kx) dx, \quad \text{with } k = 0, 1, 2, 3, \dots, \infty \quad (3.4.322.)$$

Similarly  multiplying by $\sin(kx)$ and integrating

$$b_k = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(kx) dx \quad \text{with } k = 1, 2, 3, \dots, \infty \quad (3.4.323.)$$

Coefficients so determined are called Fourier coefficients

Associated Fourier series is given by

$$f(x) = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx + \frac{1}{\pi} \sum_{n=1}^{\infty} \int_{-\pi}^{\pi} dx' f(x') [\cos(nx) \cos(nx') + \sin(nx) \sin(nx')] \quad (3.4.324.)$$

which can be written in a more compact form as

$$f(x) = \langle f \rangle + \frac{1}{\pi} \sum_{n=1}^{\infty} \int_{-\pi}^{\pi} dx' f(x') \cos[n(x - x')] \quad (3.4.325.)$$

The interval $[-\pi, \pi]$ was chosen rather arbitrarily

later on \rightarrow we will consider other intervals

Corollary 3.4.1.

Since $\frac{\cos(nx)}{i \sin(nx)} = \frac{1}{2} (e^{inx} \pm e^{-inx})$ we can rewrite (3.4.319.) as a power series in e^{ix}

$$f(x) = c_0 + \sum_{n=1}^{\infty} c_n e^{inx} + c_n^* e^{-inx} = \lim_{m \rightarrow \infty} \sum_{n=-m}^m c_n e^{inx} \quad (3.4.326.)$$

$$\text{with } c_n = c_{-n}^* = \frac{1}{2} (a_n - ib_n) = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) e^{-inx} dx \quad (3.4.327.)$$

Definition 3.4.2.

$f(x)$ is said to be piecewise continuous on an interval $[a, b]$ if it is defined and continuous

except possibly at a finite number of points x_k such that at each point of discontinuity left- and right-hand limits

$$f(x_k^-) = \lim_{x \rightarrow x_k^-} f(x), \quad f(x_k^+) = \lim_{x \rightarrow x_k^+} f(x) \quad (3.4.328.)$$

exist

At endpoints a, b only require limits $f(a^+)$ and $f(b^-)$ to exist

Note that we do not require that $f(x)$ be defined at x_k

Even if $f(x_k)$ is defined

it does not necessarily equal either left- or right-hand limit

A function $f(x)$ defined for all $x \in \mathbb{R}$ is piecewise continuous provided it is piecewise continuous on every bounded interval

Points x_k are known as jump discontinuities of $f(x)$

difference between + and - limits is magnitude of jump

Definition 3.4.3.

A function $f(x)$ is called piecewise C^1 on an interval $[a, b]$ if it is defined continuous and continuously differentiable except possibly at a finite number of points x_k

such that at each exceptional point left- and right-hand limits

$$f(x_k^-) = \lim_{x \rightarrow x_k^-} f(x), \quad f(x_k^+) = \lim_{x \rightarrow x_k^+} f(x),$$

$$f'(x_k^-) = \lim_{x \rightarrow x_k^-} f'(x), \quad f'(x_k^+) = \lim_{x \rightarrow x_k^+} f'(x)$$

exist

at endpoints we only require appropriate one-sided limits to exist

$$f(a^+), f(b^-), f'(a^+), f'(b^-)$$

For a piecewise continuous C^1 function

an exceptional point x_k is either:

(i) a jump discontinuity of f

but where left- and right-hand derivatives exist

(ii) a corner

meaning a point where f is continuous \Rightarrow so $f(x_k^-) = f(x_k^+)$

but has different left- and right-hand derivatives $f'(x_k^-) \neq f'(x_k^+)$

Lemma 3.4.1. [Riemann-Lebesgue Lemma]

If $g(x)$ is piecewise continuous in interval $[a, b]$

$$\text{then } \lim_{s \rightarrow \infty} \int_a^b g(x) \sin(sx + \alpha) dx = 0 \quad (3.4.329.)$$

Proof.

If $g(x) \in C^1([a, b])$ \Rightarrow then integration by parts leads to

$$\begin{aligned} \int_a^b g(x) \sin(sx + \alpha) dx &= - \int_a^b g(x) \frac{d}{dx} \left[\frac{\cos(sx + \alpha)}{s} \right] dx \\ &= -g(x) \frac{\cos(sx + \alpha)}{s} \Big|_a^b + \int_a^b g'(x) \frac{\cos(sx + \alpha)}{s} dx \end{aligned}$$

which goes to zero as $s \rightarrow \infty$

same reasoning is valid if $g(x)$ is differentiable

except perhaps at a finite number of points in interval

where one-sided directional derivatives must exist and be finite

Next we discuss case in which $g(x)$ is continuous in $[a, b]$

Consider a partition of closed interval $[a, b]$, $x_0 = a, x_1, \dots, x_n = b$
with $x_i - x_{i-1} = (b - a)/n$

It follows that



$$\begin{aligned} \left| \int_a^b g(x) \sin(sx + \alpha) dx \right| &= \left| \sum_{i=1}^n \int_{x_{i-1}}^{x_i} g(x) \sin(sx + \alpha) dx \right| \\ &= \left| \sum_{i=1}^n \left\{ g(x_i) \int_{x_{i-1}}^{x_i} \sin(sx + \alpha) dx \right. \right. \\ &\quad \left. \left. + \int_{x_{i-1}}^{x_i} [g(x) - g(x_i)] \sin(sx + \alpha) dx \right\} \right| \\ &\leq \sum_{i=1}^n \left[|g(x_i)| \frac{|\cos(sx_i + \alpha) - \cos(sx_{i-1} + \alpha)|}{s} \right. \\ &\quad \left. + \frac{m_i(b-a)}{n} \right] \\ &\leq \frac{2Mn}{s} + M_n (b-a) \end{aligned}$$

(3.4.330.)



where $\left\{ \begin{array}{l} M \text{ is maximum value of } g \text{ in } [a, b] \\ m_i \text{ is maximum value of } |g(x) - g(x_i)| \text{ in } [x_{i-1}, x_i] \\ M_n \text{ is maximum value of } m_i \end{array} \right.$

With this in mind,
$$\lim_{s \rightarrow \infty} \left| \int_a^b g(x) \sin(sx) dx \right| \leq M_n(b-a) \quad (3.4.329.)$$

Note that $\lim_{n \rightarrow \infty} M_n = 0$ because g is continuous and so M_n can be as small as desired by increasing n

Same applies if we replace $\sin(sx + \alpha)$ by $\cos(sx + \alpha)$

This proves lemma even if $g(x)$ is not differentiable at any point

If g is piecewise continuous

we can separate jump discontinuity points x_c

through integrals of form
$$\int_{x_c - \epsilon}^{x_c + \epsilon} g(x) \sin(sx + \alpha) dx$$

which go to zero for $\epsilon \rightarrow 0$ because g is bounded

We repeat previous reasoning in remaining intervals where g is continuous

Definition 3.4.3. ++

A function is said to have period $2L$ if $f(x + 2L) = f(x) \forall x$

For notational simplicity

we shall restrict our discussion to functions of period 2π

There is no loss of generality in doing so

since we can always use a simple change of scale $x = Ly/\pi$

to convert a function of period $2L$ into one of period 2π
the rescaled function $F(y) = f(Ly/\pi)$ lives on $[-\pi, \pi]$

Theorem 3.4.1. [Fourier convergence theorem]

If $f(x)$ is any 2π periodic piecewise C^1 function

then \Rightarrow for any $x \in \mathbb{R}$ its Fourier series converges to:

$f(x)$ if f is continuous at x

$\frac{1}{2}[f(x^+) + f(x^-)]$ if x is a jump discontinuity

(3.4.330.)

Fourier series converges to $f(x)$ at all points of continuity

At discontinuities \Rightarrow Fourier series cannot decide

whether to converge to right- or left-hand limit

and so ends up **splitting difference** by converging to their average

Proof.

Part I Let $f(x)$ be a differentiable function (therefore continuous) in interval $[a, b]$

We will show that identity (3.4.325.) is valid $\forall x \in (-\pi, \pi)$

If in addition $\rightarrow f(-\pi) = f(\pi)$

we will show that series converges to $f(x)$ in $[-\pi, \pi]$

n -th partial sum of (3.4.325.) is given by

$$S_n(x) = \frac{1}{2}a_0 + \sum_{m=1}^n a_m \cos(mx) + b_m \sin(mx)$$

$$= \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) dt \left\{ \frac{1}{2} + \sum_{m=1}^{\infty} \cos(mx) \cos(mt) + \sin(mx) \sin(mt) \right\}$$

$$= \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) \left\{ \frac{1}{2} + \sum_{m=1}^n \cos[(m(t-x))] \right\} dt$$

$$= \frac{1}{\pi} \int_{-\pi}^{\pi} f(t) K_n(t-x) dt \quad (3.4.331.)$$

$$\text{with } \rightarrow K_n(s) = \frac{1}{2} + \sum_{m=1}^n \cos(ms) \quad (3.4.332.)$$

We multiply both sides of (3.4.332.) by $\sin(s/2)$

and use relation $2 \sin \alpha \cos \beta = \sin(\alpha + \beta) + \sin(\alpha - \beta)$ to obtain

$$\sin\left(\frac{s}{2}\right) K_n(s) = \frac{1}{2} \left\{ \sin\left(\frac{s}{2}\right) + \sum_{k=1}^n \left[\sin\left(k + \frac{1}{2}\right) s \right] - \sin\left[\left(k - \frac{1}{2}\right) s \right] \right\}$$

one recognizes a telescopic sum

↳ all terms except $\sin\left[\left(n + \frac{1}{2}\right) s\right]$ cancel

Therefore

$$\text{↳ } K_n(s) = \frac{\sin\left[\left(n + \frac{1}{2}\right) s\right]}{2 \sin(s/2)} \quad (3.4.333.)$$

Note that

$$K_n(2k\pi) = \lim_{s \rightarrow 2k\pi} K_n(s) = n + \frac{1}{2} \quad k \in \mathbb{Z} \quad (3.4.334.)$$

it is easily seen from sum of cosines that defines $K_n(s)$ that 

$$\int_{-\pi}^{\pi} K_n(t - x) dt = \pi \quad (3.4.335.)$$

Therefore \Rightarrow by adding and subtracting $f(x)$ to (3.4.331.) we have

$$\begin{aligned} S_n(x) &= f(x) + \frac{1}{\pi} \int_{-\pi}^{\pi} [f(t) - f(x)] K_n(t-x) dt && (3.4.336.) \\ &= f(x) + \frac{1}{\pi} \int_{-\pi}^{\pi} \frac{f(t) - f(x)}{2 \sin[(t-x)/2]} \sin \left[\left(n + \frac{1}{2} \right) (t-x) \right] dt \end{aligned}$$

For $t \neq x \Rightarrow$ the function

$$g(t) = \frac{f(t) - f(x)}{2 \sin[(t-x)/2]} \quad (3.4.337.)$$

is continuous $\forall t \in [-\pi, \pi]$

Since $f(x) \in C^1[-\pi, \pi]$ we have

$$g(x) = \lim_{t \rightarrow x} \frac{f(t) - f(x)}{2 \sin[(t-x)/2]} = f'(x) \quad (3.4.338.)$$

and so integral in (3.4.336.) vanishes for $n \rightarrow \infty$

yielding \Rightarrow
$$\lim_{n \rightarrow \infty} S_n(x) = f(x) \quad \forall x \in (-\pi, \pi) \quad (3.4.339.)$$

If $x = \pm\pi$ \Rightarrow the denominator of $g(t)$ is canceled
for both $t \rightarrow \pi$ and $t \rightarrow -\pi$

However \Rightarrow if $x = \pm\pi$, $K_n(t - x)$ is an even function of t

$$\begin{aligned}
 K_n(t \mp \pi) &= \frac{\sin \left[\left(n + \frac{1}{2} \right) (t \mp \pi) \right]}{2 \sin[(t \mp \pi)/2]} \\
 &= \frac{\sin \left[\left(n + \frac{1}{2} \right) (-t \pm \pi) \right]}{2 \sin[(-t \pm \pi)/2]} \\
 &= \frac{\sin \left[\left(n + \frac{1}{2} \right) (-t \mp \pi) \right]}{2 \sin[(-t \mp \pi)/2]} \quad (3.4.340.)
 \end{aligned}$$

Thus $\Rightarrow \int_{-\pi}^0 K_n(t - x) dt = \int_0^{\pi} K_n(t - x) dt = \frac{1}{2} \pi \quad (3.4.341.)$

For $x = \pm\pi$ we can then write

$$S_n(x) = \frac{f(\pi) + f(-\pi)}{2} + \frac{1}{\pi} \int_0^\pi \frac{f(t) - f(\pi)}{2 \sin[(t-x)/2]} \sin \left[\left(n + \frac{1}{2}\right)(t-x) \right] dt \\ + \frac{1}{\pi} \int_{-\pi}^0 \frac{f(t) - f(-\pi)}{2 \sin[(t-x)/2]} \sin \left[\left(n + \frac{1}{2}\right)(t-x) \right] dt \quad (3.4.342.)$$

with $\rightarrow \lim_{t \rightarrow \pm\pi} \frac{f(t) - f(\pm\pi)}{2 \sin[(t-x)/2]} = f'(\pm\pi) \quad (3.4.343.)$

Applying Riemann-Lebesgue lemma to (3.4.342.) we obtain

$$\lim_{n \rightarrow \infty} S_n(\pm\pi) = \frac{f(\pi) + f(-\pi)}{2} \quad (3.4.344.)$$

which gives desired result \rightarrow if $f(\pi) = f(-\pi)$

Part II Let $f(x)$ be a 2π periodic piecewise C^1 function

We will show that (3.4.330.) holds for any $x \in \mathbb{R}$

If x is a point of differentiability of function f

we repeat previous reasoning

If x_c is corner \Rightarrow function $g(t)$ as given in (3.4.331.)

is continuous for $t \neq x$ \Rightarrow and remains bounded for $t \rightarrow x$

satisfying hypotheses of Riemann-Lebesgue Lemma

Therefore \Rightarrow Fourier series also converges to f

If function f has a jump at point x

we can use (3.3.341.) and 2π periodicity of kernel K_n to write

$$S_n(x) - \frac{f(x^+) + f(x^-)}{2} = \frac{1}{2} \int_{x-\pi}^x K_n(x-t)[f(t) - f(x^-)] dt \\ + \frac{1}{2} \int_x^{x+\pi} K_n(x-t)[f(t) - f(x^+)] dt$$

By a similar argument to one used in case

for which x is a point of differentiability of function f

we obtain desired outcome \Rightarrow that is $\lim_{n \rightarrow \infty} S_n(x) = \frac{1}{2} [f(x^+) + f(x^-)]$

Definition 3.4.5.

Most familiar convergence mechanism for sequence of functions $S_n(x)$ is pointwise convergence.

This requires that functions' values at each individual point converge in usual sense:

$$\lim_{n \rightarrow \infty} S_n(x) = f(x) \quad \forall x \in I \in \mathbb{R}e$$

Pointwise convergence requires that for every $\epsilon > 0$ and every $x \in I$ there exists an integer N depending on ϵ and x such that

$$|S_n(x) - f(x)| < \epsilon \quad \forall n \geq N$$

Pointwise convergence can be viewed as the function space version of the convergence of the components of a vector

Definition 3.4.6.

A stronger mode of convergence

is defined by demanding all points

to approach at more or less same rate to limit function

More precisely \Rightarrow a sequence of functions $S_n(x)$

is said to converge uniformly to a function $f(x)$ on a subset $I \subset \mathbb{R}$

if for any $\epsilon > 0$

there exists an integer N -- depending solely on ϵ -- such that

$$|S_n(x) - f(x)| < \epsilon \quad \forall x \in I \text{ and } \forall n \geq N \quad (3.4.348.)$$

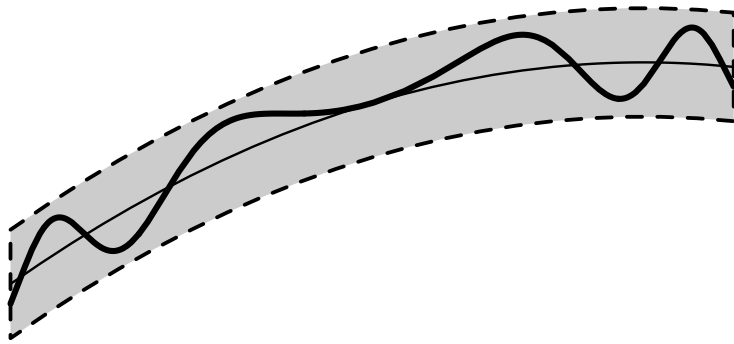
Uniformly convergent sequence of functions converges pointwise

but converse does not hold

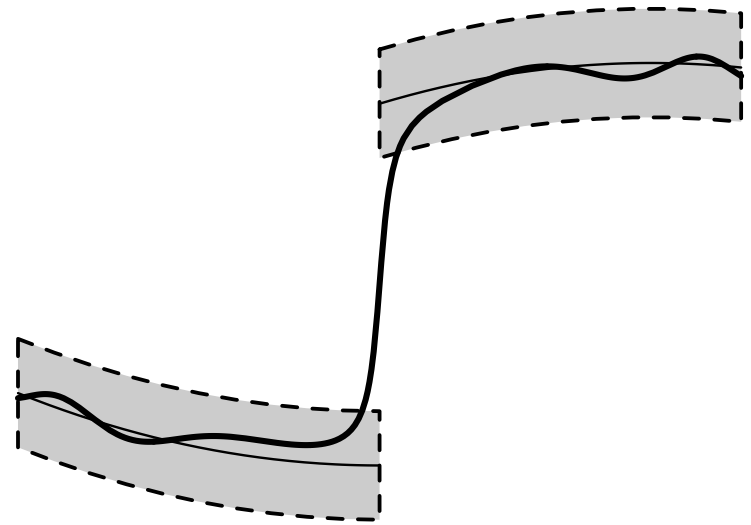
Key difference and reason for term uniform convergence

is that integer N depends only upon ϵ and not on point $x \in I$

According to (3.4.348.) the sequence converges uniformly if and only if for any small ϵ graphs of functions eventually lie inside a band of width 2ϵ centered around graph of limiting function



uniform



non-uniform

A key feature of uniform convergence is that it preserves continuity

Theorem 3.4.2.

If each $S_n(x)$ is continuous and $S_n(x) \rightarrow f(x)$ converges uniformly
then $f(x)$ is also a continuous function

Proof.

The proof is by **reductio ad absurdum**

Intuitively \Rightarrow if $f(x)$ were to have a discontinuity

then a sufficiently small band around its graph
would not connect together

and this prevents the connected graph of any continuous function
such as $S_n(x)$

from remaining entirely within the band

Uniform convergence demands

all points to behave similarly in their approaching limit

Therefore \Rightarrow it respects continuity and integration

but that mode of convergence is -- as expected -- not easy to get

For a normed vector space

one might instead consider an error of form

$$\lim_{n \rightarrow \infty} \|S_n(x) - f(x)\| \rightarrow 0 \quad (3.4.349.)$$

This last one seems to be one of the best ways

of measuring error in case of Fourier series

For finite-dimensional vector spaces such as \mathbb{R}^n

convergence in norm is equivalent to ordinary convergence

On infinite-dimensional function spaces

convergence in norm differs from pointwise convergence

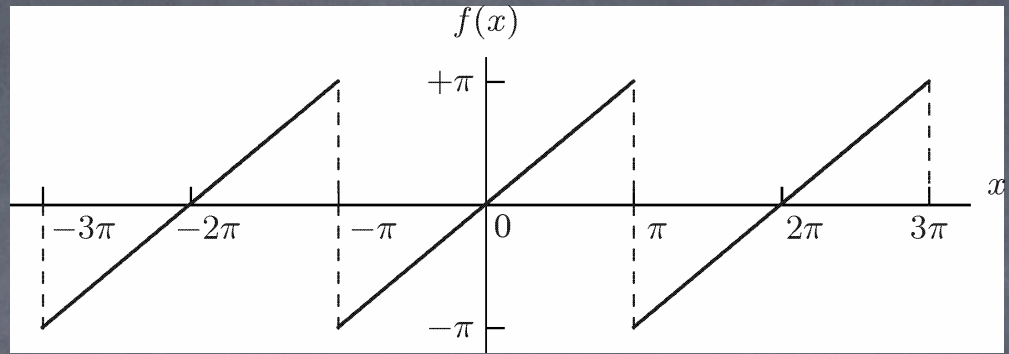
For instance \Rightarrow it is possible to construct a sequence of functions that converges in norm but does not converge pointwise **anywhere**

(We'll see this is the case in Example 3.4.3.)

Example 3.4.1.

For $f(x) = x$ with $x \in [-\pi, \pi]$

we obtain



$$a_n = 0 \forall n \quad \text{and} \quad b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} x \sin(nx) dx = \frac{2(-1)^{n+1}}{n}$$

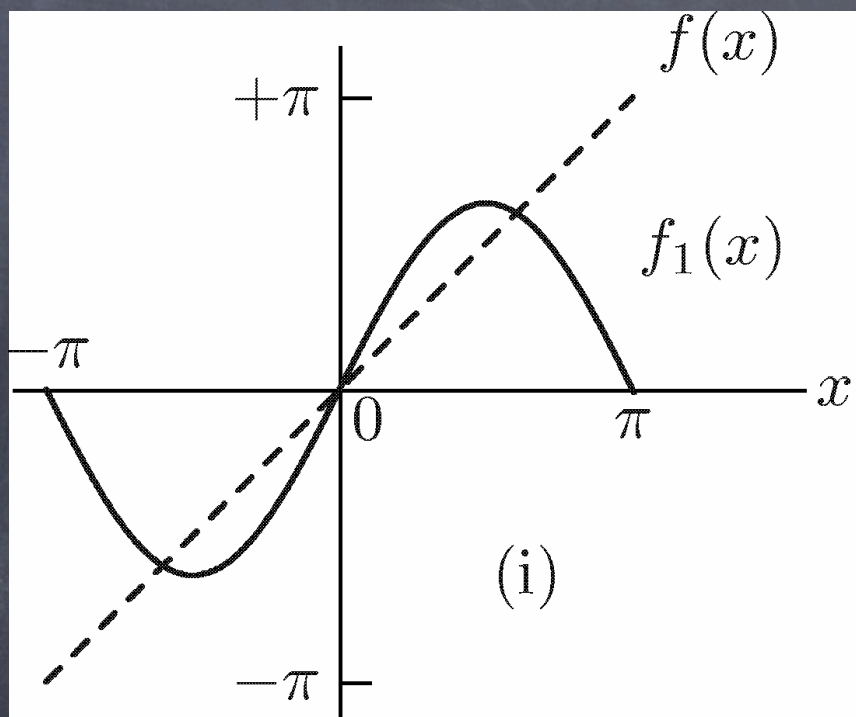
$$\text{Therefore} \quad x = 2 \sum_{n=1}^{\infty} (-1)^{n+1} \frac{\sin(nx)}{n}, \quad |x| < \pi$$

$$\text{As } x = \pm\pi \text{ series converges to } \frac{1}{2} [f(\pi) + f(-\pi)] = 0$$

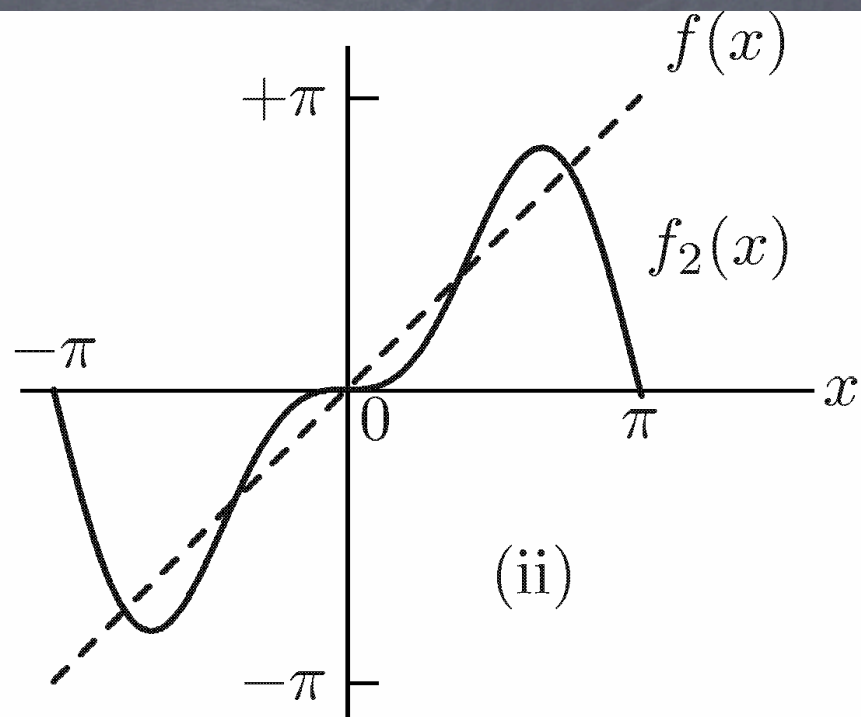
Series converges to 2π periodic extension

$$f(x) = x - 2n\pi \quad \text{if} \quad -\pi + 2n\pi < x < \pi + 2n\pi$$

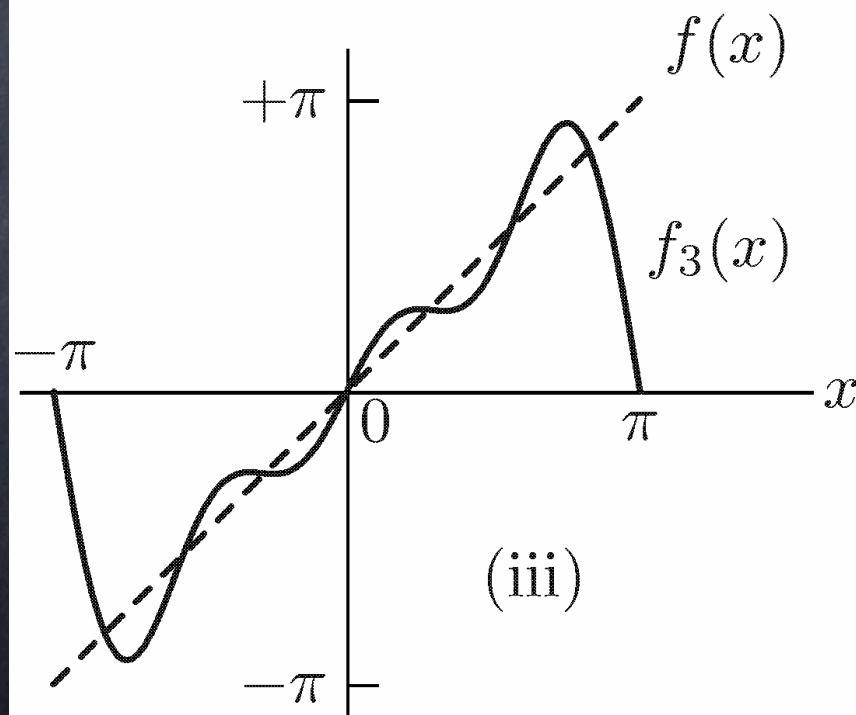
and is discontinuous at $x = \pm\pi + 2n\pi$



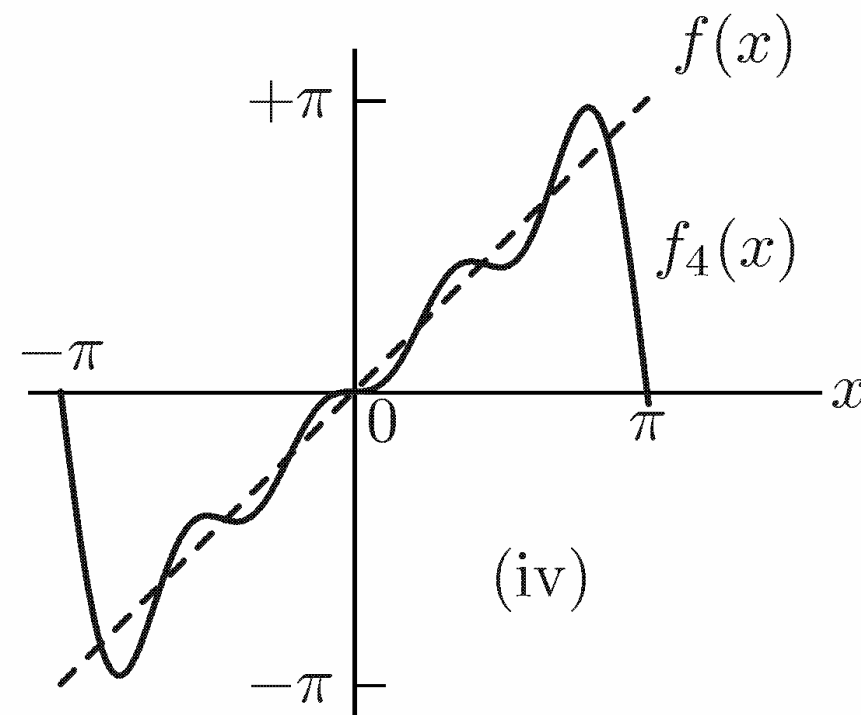
(i)



(ii)



(iii)



(iv)

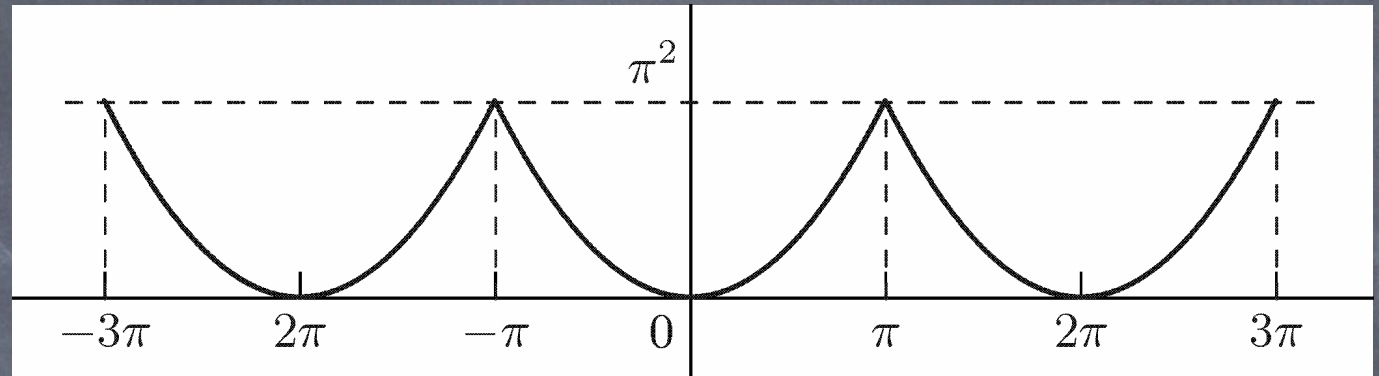
Example 3.4.2.

For $f(x) = x^2$

$$b_n = 0$$

$$a_0 = \frac{2}{3} \pi^2$$

$$a_n = \frac{4}{n^2} \cos(n\pi) \quad \text{if } n \geq 1$$



Therefore $\Rightarrow x^2 = \frac{\pi^2}{3} + 4 \sum_{n=1}^{\infty} (-1)^n \frac{\cos(nx)}{n^2} \quad |x| \leq \pi$

Fourier expansion also converges at $x = \pm\pi$ because $f(\pi) = f(-\pi)$

For $x = 0$ and $x = \pi$ Fourier expansion leads to

$$\sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n^2} = \frac{\pi^2}{12} \quad \text{and} \quad \sum_{n=1}^{\infty} \frac{1}{n^2} = \frac{\pi^2}{6}$$

Example 3.4.3.

For

$$f(x) = \begin{cases} (2a)^{-1} & \text{if } |x| < a < \pi \\ 0 & \text{if } |x| > a \end{cases}$$

we obtain $b_n = 0$, $a_n = \sin(na)/(n\pi a)$ if $n \geq 1$ and $a_0 = 1/\pi$

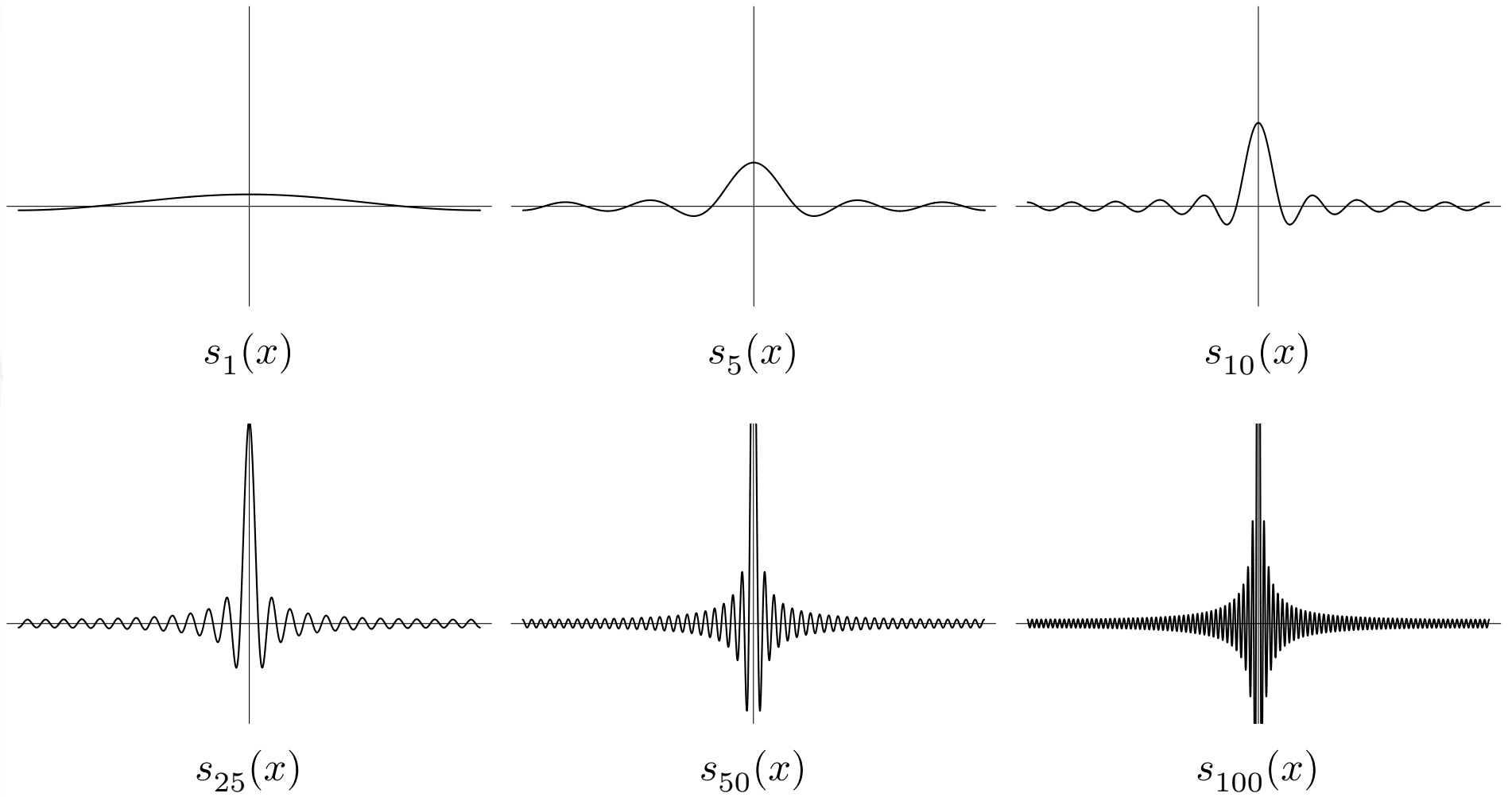
Therefore \Rightarrow

$$f(x) = \frac{1}{\pi} \left[\frac{1}{2} + \sum_{n=1}^{\infty} \frac{\sin(na)}{na} \cos(nx) \right], \quad |x| \leq \pi$$

For $x = \pm a$ series converges to $(4a)^{-1}$

If $a \rightarrow 0$ the $f(x) \rightarrow \delta(x)$ with

$$\delta(x) = \frac{1}{\pi} \left[\frac{1}{2} + \sum_{n=1}^{\infty} \cos(nx) \right] = \frac{1}{2\pi} \lim_{n \rightarrow \infty} \sum_{n=-m}^m e^{inx}, \quad |x| \leq \pi$$



We had to truncate last two graphs; spikes extend beyond the top

Series does not converge pointwise

but it converges as a distribution to $\delta(x)$ for $|x| \leq \pi$

Indeed \rightarrow n -th partial sum

$$S_n(x) = \frac{1}{\pi} \left[\frac{1}{2} + \sum_{n=1}^{\infty} \cos(nx) \right] = \frac{\sin \left[\left(n + \frac{1}{2} \right) x \right]}{2\pi \sin(x/2)}$$

satisfies $\int_{-\pi}^{\pi} S_n(x) dx = 1, \quad \forall n \geq 0$

and

$$\begin{aligned} \lim_{n \rightarrow \infty} \int_{-\pi}^{\pi} S_n(x) f(x) dx &= f(0) + \lim_{n \rightarrow \infty} \int_{-\pi}^{\pi} \frac{f(x) - f(0)}{\sin(x/2)} \sin \left[\left(n + \frac{1}{2} \right) x \right] dx \\ &= f(0) \quad \forall f \text{ test functions} \end{aligned}$$

Actually \rightarrow Fourier series converges to

2π periodic extension of original function

$$\lim_{m \rightarrow \infty} \sum_{n=-m}^m \delta(x - 2n\pi) = \frac{1}{\pi} \left[\frac{1}{2} + \sum_{n=1}^{\infty} \cos(nx) \right] = \lim_{m \rightarrow \infty} \frac{1}{2\pi} \sum_{n=-m}^m e^{inx}$$

consisting of a periodic array of delta spikes

concentrated at all integer multiples of 2π

Example 3.4.4.

For

$$f(x) = \begin{cases} 1 & \text{if } |x| < 1/2 \\ 0 & \text{if } |x| > 1/2 \end{cases} \quad x \in [-1, 1] \quad (3.4.357.)$$

we obtain $b_n = \int_{-1}^1 f(x) \sin(n\pi x) dx = 0$ $a_0 = 1$

and $a_n = \int_{-1}^1 f(x) \cos(n\pi x) dx = \begin{cases} (-1)^{(n+1)/2} 2/(n\pi) & n \text{ odd} \\ 0 & n \text{ even} \end{cases}$

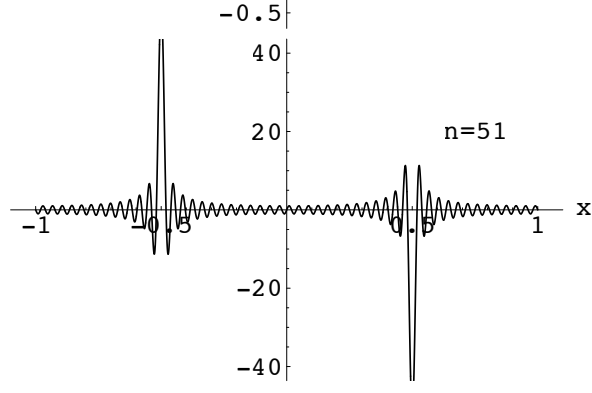
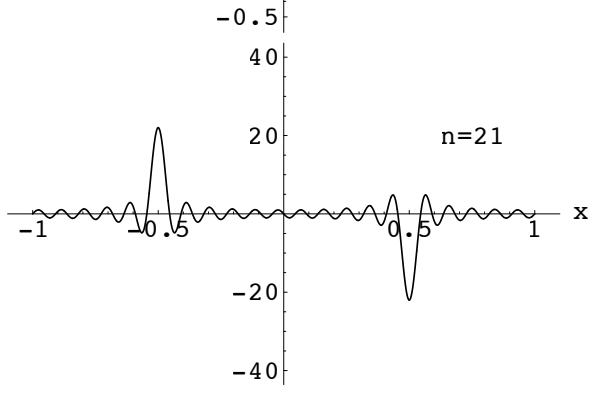
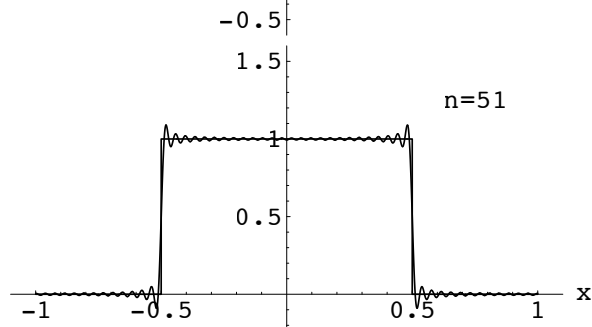
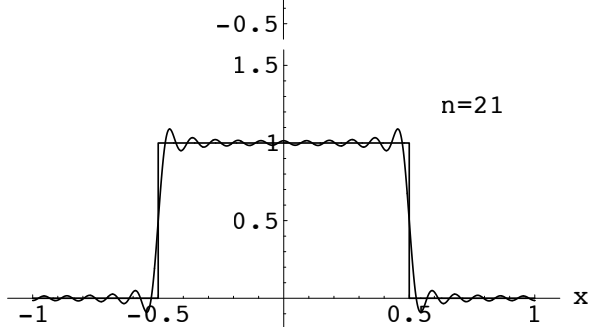
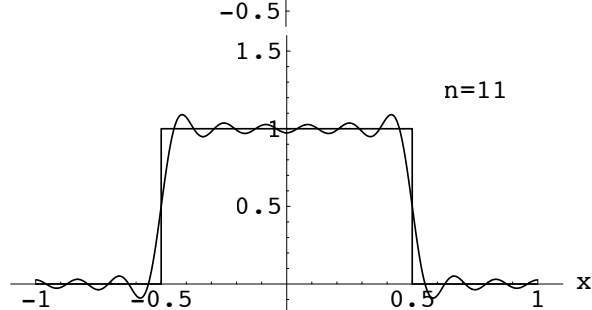
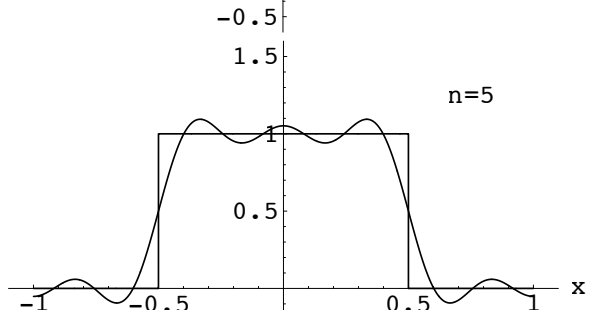
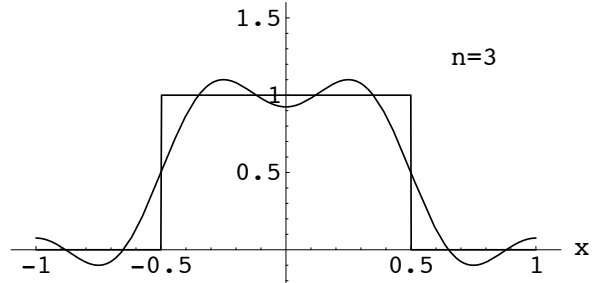
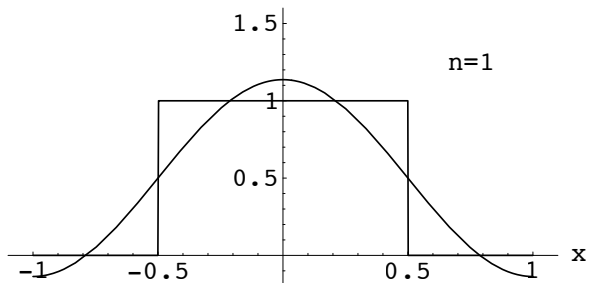
n -th partial sum reads $S_n(x) = \frac{1}{2}a_0 + \sum_{m=1}^n a_m \cos(m\pi x)$

$$dS_n(x)/dx \rightarrow \delta\left(x + \frac{1}{2}\right) - \delta\left(x - \frac{1}{2}\right)$$

Near jumps there should be a consistent overshoot of about 9%

So-called **Gibbs overshoot** is a manifestation

of subtle non-uniform convergence of Fourier series



Definition 3.4.9.

Consider complex form of Fourier series of $f(x) : [-L, L] \rightarrow \mathbb{R}$

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{in\pi x/L}, \quad c_n = \frac{1}{2L} \int_{-L}^L f(x) e^{-in\pi x/L} dx$$

We can rewrite this series as

$$f(x) = \frac{1}{\sqrt{2\pi}} \sum_k \hat{f}(k) e^{ikx} \Delta k \quad (3.4.370.)$$

where $k = n\pi/L$, $\Delta k = \pi/L$

$$\text{and } \hat{f}(k) = \sqrt{2\pi} c_n \frac{L}{\pi} = \frac{1}{\sqrt{2\pi}} \int_{-L}^L f(x) e^{ikx} dx \quad (3.4.373.)$$

Let us now consider limit $L \rightarrow \infty$

In such a case $\Delta k \rightarrow 0$ while (3.4.370.) and (3.4.373.)

$$\hat{f}(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx \quad (3.4.375.)$$

approach \Rightarrow

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(k) e^{ikx} dk \quad (3.4.379.)$$

assuming that both integrals converge

$\hat{f} : \mathbb{R} \rightarrow \mathbb{R}$ **Fourier transform** of function $f (f : \mathbb{R} \rightarrow \mathbb{R})$

Fourier transform operator $\hat{f}(k) = \mathcal{F}[f(x)]$

maps each (sufficiently nice) function of spatial variable x
to a function of frequency variable k

Expression which retrieves f from \hat{f}

is inverse of Fourier transform operator $f(x) = \mathcal{F}^{-1} \left[\hat{f}(k) \right]$

\mathcal{F} & \mathcal{F}^{-1} are generalization of Fourier series

for functions f defined on $(-\infty, \infty)$

i.e. $f(x) \in \mathcal{L}^2$ is square-integrable

Before proceeding we show validity of (3.4.375.) and (3.4.379.)

for functions f which are piecewise C^1 and satisfy $\int_{-\infty}^{\infty} |f(x)| dx < \infty$

Proof is very similar to one carried out for Fourier series

(3.4.375.) and (3.4.379.) entail

$$f(x) = \int_{-\infty}^{\infty} \frac{1}{2\pi} \left[\int_{-\infty}^{\infty} e^{ik(x-x')} dk \right] f(x') dx'$$

so what must be shown is that

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ik(x-x')} dk = \delta(x-x') \quad \rightarrow \quad \int_{-\infty}^{\infty} = \lim_{r \rightarrow \infty} \int_{-r}^r \quad (3.4.382.)$$

$$\text{Indeed } \rightarrow \quad \frac{1}{2\pi} \int_{-r}^r e^{ikt} dk = \frac{1}{2\pi} \frac{e^{irt} - e^{-irt}}{it} = \frac{1}{\pi} \frac{\sin(rt)}{t}$$

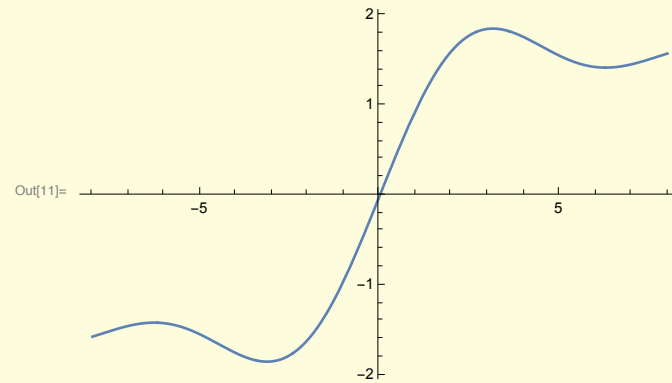
$$\text{with } \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{\sin(rt)}{t} dt = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{\sin(u)}{u} du = 1$$



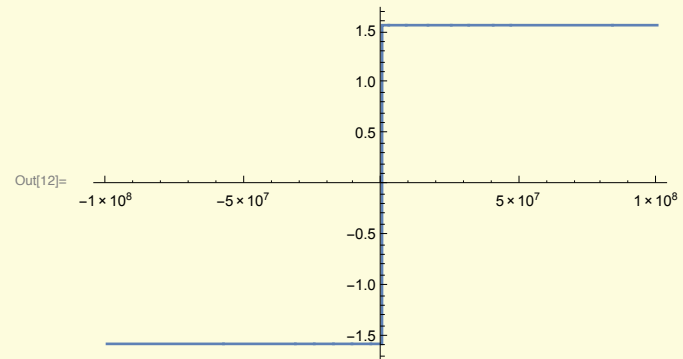
```
In[6]:= Integrate[(Sin[x]) / x, x]
```

```
Out[6]= SinIntegral[x]
```

```
In[11]:= Plot[SinIntegral[x], {x, -8, 8}]
```



```
In[12]:= Plot[SinIntegral[x], {x, -100 000 000, 100 000 000}]
```



Therefore



$$\begin{aligned} f(x) &= \lim_{r \rightarrow \infty} \int_{-\infty}^{\infty} \left[\frac{1}{2\pi} \int_{-r}^r e^{ik(x-x')} dk \right] f(x') dx' \\ &= \lim_{r \rightarrow \infty} \int_{-\infty}^{\infty} \frac{\sin[r(x-x')]}{\pi(x-x')} f(x') dx' \\ &= \lim_{r \rightarrow \infty} \int_{-\infty}^{\infty} \frac{\sin(rt)}{\pi t} [f(x+t) - f(x) + f(x)] dt \\ &= f(x) + \lim_{r \rightarrow \infty} \int_{-\infty}^{\infty} \sin(rt) \frac{f(x+t) - f(x)}{\pi t} dt \end{aligned}$$

For $r \rightarrow \infty$ second term cancels

because $[f(x+t) - f(x)]/t$ remains bounded for $t \rightarrow 0$

Equation (3.4.382.) also implies

$$\int_{-\infty}^{\infty} \varphi_{k'}^*(x) \varphi_k(x) dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ix(k-k')} dx = \delta(k - k') \quad (3.4.386.)$$

indicating that functions $\varphi_k(x) = e^{ikx} / \sqrt{2\pi}$ are orthogonal with

with respect to inner product $\langle u, v \rangle = \int_{-\infty}^{\infty} u^*(x) v(x) dx$

and are **normalized** with respect to variable

Note that convergence of integrals (3.4.382.) and (3.4.386.)
should be understood as distributions

Example 3.4.5.

Fourier transform of rectangular pulse,

$$f(x) = \Theta(x+a) - \Theta(x-a) = \begin{cases} 1 & -a < x < a \\ 0 & |x| > a \end{cases} \quad (3.4.389.)$$

(a.k.a. box function of width $2a$) is easily computed

$$\hat{f}(k) = \frac{1}{\sqrt{2\pi}} \int_{-a}^a e^{-ikx} dx = \frac{e^{ika} - e^{-ika}}{\sqrt{2\pi}ik} = \sqrt{\frac{2}{\pi}} \frac{\sin(ak)}{k} \quad (3.4.390.)$$

Reconstruction of pulse via inverse transform (3.4.379.) tells us that

$$\frac{1}{\pi} \int_{-\infty}^{\infty} \frac{e^{ikx} \sin(ak)}{k} dk = \begin{cases} 1 & -a < x < a \\ 1/2 & x = \pm a \\ 0 & |x| > a \end{cases} \quad (3.4.391.)$$

Note convergence to middle of jump discontinuities at $x = \pm a$

Real part of this complex integral
produces a striking trigonometric integral identity

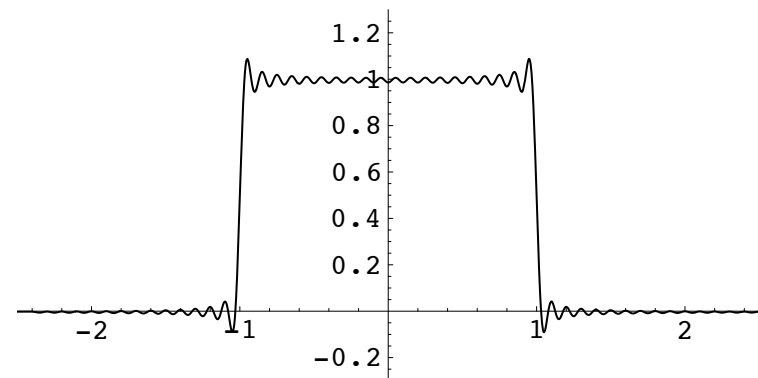
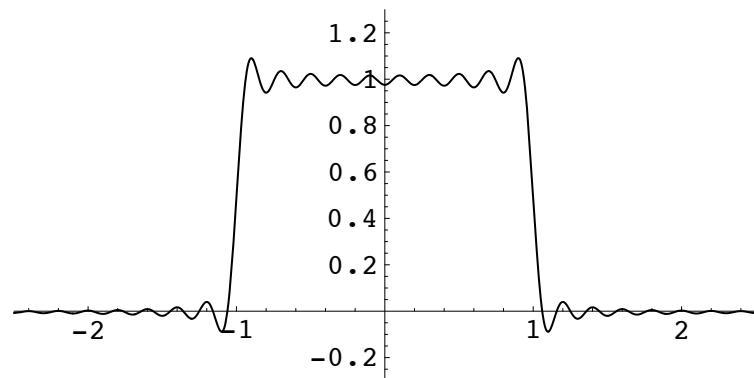
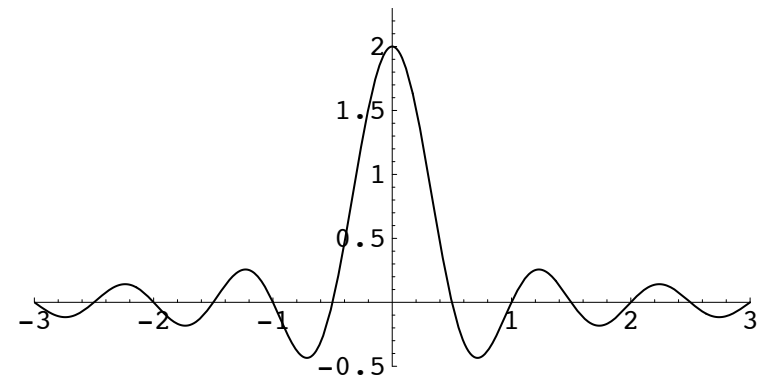
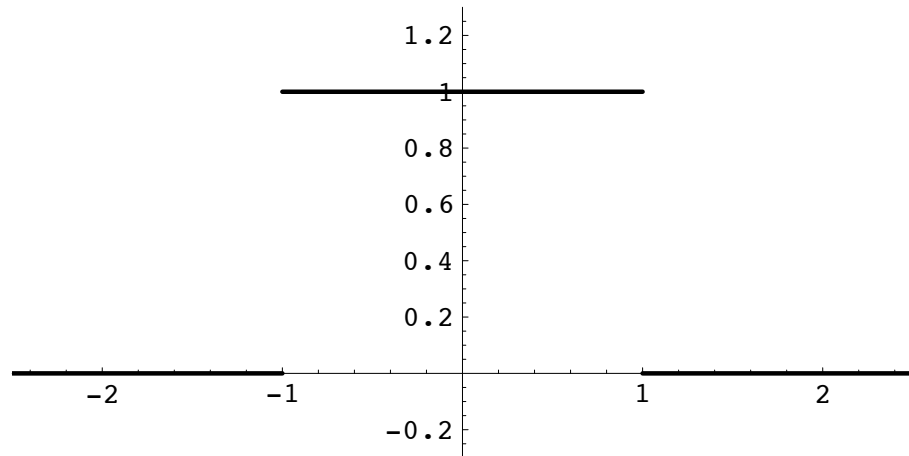
$$\frac{1}{\pi} \int_{-\infty}^{\infty} \frac{\cos(kx) \sin(ak)}{k} dk = \begin{cases} 1 & -a < x < a \\ 1/2 & x = \pm a \\ 0 & |x| > a \end{cases}$$

identity resulting from imaginary part

$$\frac{1}{\pi} \int_{-\infty}^{\infty} \frac{\sin(kx) \sin(ak)}{k} dk = 0$$

is not surprising because integrand is odd

Non-uniform convergence of the integral leads to appearance of a Gibbs phenomenon at two discontinuities similar to the non-uniform convergence of a Fourier series



Since we are dealing with an infinite integral must break off numerical integrator restricting it to a finite interval

Left graph is obtained by integrating from $-5 \leq k \leq 5$

Right graph is obtained by integrating from $-10 \leq k \leq 10$

